

On certain infinite integrals involving Struve functions and parabolic cylinder functions

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The object of the present note is to obtain a number of infinite integrals involving Struve functions and parabolic cylinder functions.

1. G. N. Watson⁽¹⁾ has proved that

$$\int_0^\infty e^{-\frac{1}{2}x^2} x^m D_n(x) dx = \frac{\sqrt{\pi} 2^{\frac{1}{2}(n-m-1)} \Gamma(m+1)}{\Gamma(\frac{1}{2}m - \frac{1}{2}n + 1)} \quad (R(m) > -1). \quad (1)$$

From (1)

$$\begin{aligned} & \int_0^\infty e^{-\frac{1}{2}x^2} x^m {}_pF_q(a_1, \dots, a_p; c_1, \dots, c_q; x^2 y^2) D_n(x) dx \\ &= \frac{\sqrt{\pi} 2^{\frac{1}{2}(n-m-1)} \Gamma(m+1)}{\Gamma(\frac{1}{2}m - \frac{1}{2}n + 1)} {}_{p+2}F_{q+1} \left[a_1, \dots, a_p, \frac{1}{2}m + \frac{1}{2}, \frac{1}{2}m + 1; \right. \\ & \left. c_1, \dots, c_q, \frac{1}{2}m - \frac{1}{2}n + 1; 2y^2 \right] \quad (2) \end{aligned}$$

follows provided that the integral is convergent and term-by-term integration is permissible. A great many interesting particular cases of (2) are easily deducible: the following will be used in this paper.

$$\int_0^\infty \sqrt{xy} x^s e^{-\frac{1}{2}x^2} \mathbf{H}_{s+\frac{1}{2}}(xy) D_{2s+2}(x) dx = (-1)^s y^{s+\frac{1}{2}} e^{-\frac{1}{2}y^2} D_{2s+1}(y) \quad (3)$$

$$\int_0^\infty \sqrt{xy} x^{s+\frac{1}{2}} e^{-\frac{1}{2}x^2} J_s(xy) D_{2s+2}(x) dx = (-1)^{s+\frac{1}{2}} y^{s+\frac{1}{2}} e^{-\frac{1}{2}y^2} D_{2s+2}(y) \quad (4)$$

$$\int_0^\infty e^{-\frac{1}{2}x^2} D_{2s}(x) \cos xy dx = \sqrt{\frac{\pi}{2}} (-1)^s y^{2s} e^{-\frac{1}{2}y^2} \quad (5)$$

(4) is already known.⁽²⁾ Here and later s is a non-negative integer.

2. The following integrals are obtained by integration with respect to y from 0 to ∞ and inversion of the order of integrations.

Let us divide both sides of (5) by $(y^2 + z^2)^s$, integrate, and use

$$\int_0^\infty t^{2s} (t^2 + a^2)^{-s} e^{-\frac{1}{2}t^2} dt = 2^{s-\frac{1}{2}} \Gamma(s + \frac{1}{2}) e^{\frac{1}{2}a^2} D_{-2s}(a). \quad (6)$$

We get for $s > 0$

$$\int_0^\infty x^{s-\frac{1}{2}} e^{-\frac{1}{2}x^2} K_{s-\frac{1}{2}}(xz) D_{2s}(x) dx = \sqrt{\frac{\pi}{2}} (-1)^s \Gamma(2s) z^{s-\frac{1}{2}} e^{\frac{1}{2}z^2} D_{-2s}(z). \quad (7)$$

Multiplying (4) by $K_{s+\frac{1}{2}}(yz)$ and integrating

$$\int_0^\infty x^{2s+1} (x^2 + z^2)^{-\frac{1}{2}} e^{-ix^2} D_{2s+2}(x) dx = \Gamma(2s + 2) z^{2s+1} e^{iz^2} D_{-2s-2}(z). \quad (8)$$

Multiplying (5) by e^{-vz} , integrating, and using a well-known integral representation⁽³⁾ of parabolic cylinder functions we obtain⁽¹⁾

$$\int_0^\infty z (x^2 + z^2)^{-1} e^{-ix^2} D_{2s}(x) dx = \sqrt{\frac{\pi}{2}} \Gamma(2s + 1) (-1)^s e^{iz^2} D_{-2s-1}(z). \quad (9)$$

From this formula it follows that

$$\int_0^\infty x^{2s} (x^2 + z^2)^{-1} e^{-ix^2} D_{2s}(x) dx = \sqrt{\frac{\pi}{2}} \Gamma(2s + 1) z^{2s-1} e^{iz^2} D_{-2s-1}(z), \quad (10)$$

since $x^{2s}/(x^2 + z^2) = (-)^s z^{2s}/(x^2 + z^2) +$ an even polynomial of degree $2s - 2$ in x , and the contribution of that polynomial vanishes on account of the orthogonal property of parabolic cylinder functions.

Multiplying (3) by $y K_{s+\frac{1}{2}}(yz)$, integrating, and using (10) and the known result⁽⁴⁾

$$\int_0^\infty x \mathbf{H}_{s+\frac{1}{2}}(xz) K_{s+\frac{1}{2}}(xy) dx = z^{s+\frac{1}{2}} y^{-s-\frac{1}{2}} (y^2 + z^2)^{-1},$$

we obtain

$$\int_0^\infty x^{s+\frac{1}{2}} e^{-ix^2} K_{s+\frac{1}{2}}(xy) D_{2s+1}(x) dx = \sqrt{\frac{\pi}{2}} (-1)^s \Gamma(2s + 3) y^{s-\frac{1}{2}} e^{iy^2} D_{-2s-3}(y). \quad (11)$$

3. To conclude this paper, a few integrals will be evaluated with the help of the operational calculus.

We write $f(p) \doteq h(t)$

if $f(p) = p \int_0^\infty e^{-pt} h(t) dt.$

The following results are well known⁽⁵⁾:—

$$f\left(\frac{p}{x}\right) \doteq h(xt) \quad (12) \qquad p(p^2 + x^2)^{-\frac{1}{2}} \doteq J_0(xt) \quad (13)$$

$$p I_s(y\sqrt{p}) K_s(x\sqrt{p}) \doteq \frac{1}{2t} \exp\left(-\frac{x^2 + y^2}{4t}\right) I_s\left(\frac{xy}{2t}\right) \quad (14)$$

$$\exp\left(-\frac{x^2 + y^2}{4p}\right) I_s\left(\frac{xy}{2p}\right) \doteq J_s(x\sqrt{t}) J_s(y\sqrt{t}) \quad (15)$$

$$\Gamma(s) p^{s+1} e^{ip^2} D_{-s}(p) \doteq \frac{d^s}{dt^s} (t^{s-1} e^{-it^2}), \quad (s > 0). \quad (16)$$

Goldstein⁽⁶⁾ has proved that if $\phi(p) \doteq f(t)$ and $\psi(p) \doteq g(t)$ then

$$\int_0^\infty \phi(t) g(t) t^{-1} dt = \int_0^\infty f(t) \psi(t) t^{-1} dt. \quad (17)$$

In (8) let us put $z = p$, multiply by p , and interpret by means of (13) and (16), thus obtaining

$$\int_0^\infty x^{2s+1} e^{-4x^2} J_0(xz) D_{2s+2}(x) dx = \frac{d^{2s+1}}{dz^{2s+1}} (z^{2s+1} e^{-\frac{1}{2}z^2}). \quad (18)$$

Apply (17) to (14) and (15) to find

$$\begin{aligned} & \int_0^\infty I_s(z\sqrt{t}) J_s(y\sqrt{t}) J_s(z\sqrt{t}) K_s(y\sqrt{t}) dt \\ &= \frac{1}{2} \int_0^\infty e^{-t(w^2+z^2)t} I_s^2(\frac{1}{2}yzt) dt = \frac{(-1)^{-s-\frac{1}{2}}}{\pi yz} Q_{s-\frac{1}{2}}\left(-\frac{y^4+z^4}{2y^2z^2}\right) \end{aligned} \quad (19)$$

$(y \pm z > 0).$

In the integral

$$\int_0^\infty x(x^4 + 4k^4)^{-\frac{1}{2}} J_0(xy) dx = J_0(ky) K_0(ky)$$

we put $k = \sqrt{\frac{1}{2}p}$, multiply by p , and have on interpretation

$$p J_0(y\sqrt{\frac{1}{2}p}) K_0(y\sqrt{\frac{1}{2}p}) \doteq (2t)^{-1} J_0\left(\frac{y^2}{4t}\right). \quad (20)$$

Combining this with $p/(p + 1) \doteq e^{-t}$ in the manner of (17),

$$\int_0^\infty x e^{-2x^2} J_0(xy) K_0(xy) dx = \frac{\pi}{16} \{H_0(\frac{1}{4}y^2) - Y_0(\frac{1}{4}y^2)\} \quad (21)$$

can be derived.

From (15) and (20) we deduce that

$$\begin{aligned} & \int_0^\infty J_0(x\sqrt{t}) J_0(y\sqrt{t}) J_0(z\sqrt{t}) K_0(z\sqrt{t}) dt \\ &= \frac{1}{2} \int_0^\infty e^{-t(x^2+y^2)t} J_0(\frac{1}{2}z^2t) I_0(\frac{1}{2}xyt) dt = \pi^{-1} (xyz^2i)^{-\frac{1}{2}} Q_{-\frac{1}{2}}\left(\frac{x^4+y^4-2x^2y^2+4z^4}{8xyz^2i}\right). \end{aligned} \quad (22)$$

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