FUNDAMENTAL SOLUTIONS AND SURFACE DISTRIBUTIONS

by R. A. ADAMS and G. F. ROACH¹ (Received 22nd January 1968)

Introduction

When studying the solutions of elliptic boundary value problems in a bounded, smoothly bounded domain $D \subset R_n$ we often encounter the formula

$$\int_{\partial D} \left\{ u(x) \frac{\partial}{\partial n_x} \gamma(x, y) - \gamma(x, y) \frac{\partial}{\partial n_x} u(x) \right\} dS_x = \begin{cases} u(y) & \text{if } y \in D \\ \frac{1}{2}u(y) & \text{if } y \in \partial D \\ 0 & \text{if } y \neq D \end{cases}$$
(1a)

here
$$\psi(x) \in C^2(D) \cap C'(\overline{D})$$
 is a solution of the second order self-adjoint elliptic

where $u(x) \in C^2(D) \cap C'(D)$ is a solution of the second order self-adjoint elliptic equation

$$Lu(x) \equiv (\Delta \pm k^2)u(x) = 0, \quad x \in D$$
⁽²⁾

and $\frac{\partial}{\partial n_x}$ denotes differentiation along the inward normal to ∂D at $x \in \partial D$. $\gamma(x, y)$ is a fundamental solution of (2), and as such has at x = y a singularity

 $\gamma(x, y)$ is a fundamental solution of (2), and as such has at x = y a singularity described by

$$\gamma(x, y) \sim \begin{cases} -\frac{1}{2\pi} \log |x - y| & \text{if } n = 2, \\ \frac{1}{2n - 4} \Gamma\left(\frac{n}{2}\right) \pi^{-n/2} |x - y|^{2-n} & \text{if } n \ge 3. \end{cases}$$
(3)

The results (1a) and (1c) can be obtained in a straightforward way by applying Green's Theorem to u(x) and any fundamental solution which is defined in a sufficiently large domain. However the result (1b) is neither as obvious nor as easily obtained as is generally claimed in textbooks, though, as we shall see, it is true in the sense of the theory of distributions for fundamental solutions which, apart from a singularity of type (3) at x = y, are regular in $\overline{D} \times \overline{D}$. For other choices of the fundamental solution (e.g. the Dirichlet Green's function) not satisfying this restriction, (1b) is meaningless unless a suitable definition can be given for the left hand side. In this paper we shall establish (1b) for fundamental solutions having the required behaviour in $\overline{D} \times \overline{D}$ and shall show that when a maximum principle is available ($L = \Delta - k^2$, $k^2 \ge 0$) (1b) can be made meaningful in the distributional sense for the Dirichlet Green's function of L and D. It is sufficient to demonstrate this for $L = \Delta$.

¹ Now at the Department of Mathematics, University of Strathclyde.

1. For convenience we restrict ourselves to the case n = 2. Results for $n \ge 3$ follow in a similar manner. Fix $y \in \partial D$ and let K_{ε} be the disc of radius ε centred at y. Let $S_{\varepsilon} = \partial K_{\varepsilon} \cap D$. Applying Green's Theorem to u(x) and y(x, y) in $D - K_{\varepsilon}$ and performing a simple residue calculation, noting that on S_{ε} we have

$$\gamma(x, y) \sim -\frac{1}{2\pi} \log \varepsilon, \quad \frac{\partial}{\partial n_x} = \frac{\partial}{\partial \varepsilon}$$

and $dS_x = \varepsilon d\theta \ (0 \le \theta \le \pi)$, we obtain the result

$$\lim_{\varepsilon \to 0} \int_{\partial D - K_{\varepsilon}} \left\{ u(x) \frac{\partial}{\partial n_{x}} \gamma(x, y) - \gamma(x, y) \frac{\partial}{\partial n_{x}} u(x) \right\} dS_{x} = \frac{1}{2} u(y).$$

Therefore to establish (1b) it suffices to show that

$$\lim_{\varepsilon \to 0} \int_{\partial D \cap K_{\varepsilon}} \left\{ \frac{\partial}{\partial n_{x}} \gamma(x, y) - \gamma(x, y) \frac{\partial}{\partial n_{x}} u(x) \right\} dS_{x} = 0.$$
 (4)

Since ∂D is smooth, we have in a neighbourhood of y, $dS_x \cong dr$ where r = |x-y|. Since $\frac{\partial}{\partial n_x} u(x)$ is continuous in \overline{D} it follows that

$$\lim_{\varepsilon \to 0} \int_{\partial D \cap K_{\varepsilon}} \gamma(x, y) \frac{\partial}{\partial n_{x}} u(x) dS_{x} = \lim_{\varepsilon \to 0} \frac{-1}{\pi} \frac{\partial}{\partial n_{y}} u(y) \int_{0}^{\varepsilon} \log r dr = 0.$$
(5)

Finally, let r(x) denote the radius of the circle C_x through x and y which is tangent to ∂D at x. Since ∂D is smooth, $r(x) \to R$, the radius of the osculating circle for ∂D at y, as $x \to y$. Thus for ε small enough and $x \in \partial D \cap K_{\varepsilon}$ it follows that $r(x) \ge \frac{1}{2}R$. If α is the angle between the vector from y to x and the inward normal n_x to ∂D at x we have, since $|n_x| = 1$ and $|\nabla_x|x-y|| = 1$, that

$$\left|\frac{\partial}{\partial n_x}\gamma(x, y)\right| \sim \frac{1}{2\pi} \left| n_x \cdot \nabla_x \log \left| x - y \right| \right| = \frac{1}{2\pi \left| x - y \right|} \left| \cos \alpha \right|.$$

By the geometry of the circle C_x it follows that

$$\frac{\left|\cos\alpha\right|}{\left|x-y\right|} = \frac{1}{2r(x)} \le \frac{1}{R}$$

and hence we have at once that

$$\lim_{\varepsilon \to 0} \int_{\partial D \cap K_{\varepsilon}} u(x) \frac{\partial}{\partial n_{x}} \gamma(x, y) dS_{x} = 0.$$

This completes the proof of (4) and so of (1b).

2. The harmonic Green's function, G(x, y), of D is defined for $x \in \overline{D}$, $y \notin \partial D$ by $G(x, y) = \gamma(x, y) + w(x, y)$, where $\gamma(x, y)$ is the function on the right hand side of (3) and

$$\begin{aligned} \Delta_x w(x, y) &= 0, & x \in D, \\ w(x, y) &= -\gamma(x, y), & x \in \partial D. \end{aligned}$$

G(x, y) is positive for $x, y \in D$ (1, p. 262). Since G(x, y) is not properly defined for $y \in \partial D$, the appropriate form of (1b), namely

$$\int_{\partial D} u(x) \frac{\partial}{\partial n_x} G(x, y) dS_x = \frac{1}{2} u(y), \quad y \in \partial D$$
(6)

requires interpretation.

To this end let $G_{\epsilon}(x, y)$ be the harmonic Green's function for the region $D_{\epsilon} = D \cap K_{\epsilon}$. Applying (1b) over D with $\gamma(x, y) = G_{\epsilon}(x, y)$, we obtain

$$\int_{\partial D} u(x) \frac{\partial}{\partial n_x} G_{\varepsilon}(x, y) dS_x - \int_{\partial D \cap K_{\varepsilon}} G_{\varepsilon}(x, y) \frac{\partial}{\partial n_x} u(x) dS_x = \frac{1}{2} u(y).$$

Fix ε_0 . For $\varepsilon < \varepsilon_0$, $G_{\varepsilon_0}(x, y) - G_{\varepsilon}(x, y)$ is harmonic in D_{ε} and non-negative on ∂D_{ε} . By the maximum principle it is non-negative in D_{ε} .

Thus

$$\left| \int_{\partial D \cap K_{\varepsilon}} G_{\varepsilon}(x, y) \frac{\partial}{\partial n_{x}} u(x) dS_{x} \right| \leq \text{const.} \int_{\partial D \cap K_{\varepsilon}} G_{\varepsilon_{0}}(x, y) dS_{x}.$$

As in (5) above the right hand side tends to zero with ε . Hence

$$\lim_{\varepsilon \to 0} \int_{\partial D} u(x) \frac{\partial}{\partial n_x} G_{\varepsilon}(x, y) dS_x = \frac{1}{2} u(y), \quad y \in \partial D.$$

This shows that in (6) we may interpret $\frac{\partial}{\partial n_x} G(x, y)$ as the limit in the distribution sense as $\varepsilon \to 0$ of the well defined functions $\frac{\partial}{\partial n_x} G_{\varepsilon}(x, y)$ (2, Chapter 2; 3).

REFERENCES

(1) R. COURANT and D. HILBERT, *Methods of Mathematical Physics*, Vol. II (Interscience, New York, 1962).

(2) I. M. GEL'FAND and G. E. SHILOV, *Generalised Functions*, Vol. I (Academic Press, 1964).

(3) L. SCHWARTZ, Théorie des distributions (Hermann et Cié, Paris, 1950).

DEPARTMENT OF MATHEMATICS THE UNIVERSITY OF BRITISH COLUMBIA