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A NEW METHOD OF ANALYSIS AND REDUCTION OF VLBI OBSERVATIONS

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<u>ABSTRACT</u> In this paper we develope an inertial system of coordinates introducing a new observable: the arc between a pair of radio sources. We shall discuss how the conditional equation involving all the parameters is written in different systems when the observations are performed with a two and a three-element interferometers.

INTRODUCTION

In the classical algorithm of analysis of VLBI observations, the equatorial coordinates of the radio sources, the coordinates of the stations, the Earth Orientation Parameters (EOP) and, eventually, the corrections to the position of the celestial pole given by the IAU 1976/1980 models are simultaneously adjusted.

We have explored a new method of constructing a celestial frame independently of the terrestrial frame and of the Earth's orientation in space. The new strategy of observation consists on "ideally" observing two extragalactic radio sources simultaneously with two pairs of radiotelescopes, each one at the extremities of a VLBI baseline. On the basis of the classical VLBI observables -the delay and the delay rate- we have constructed a new observable: the arc between a pair of sources. A set of these new observables leads to spherical coordinates of the radio sources relative to an arbitrary plane and origin.

We have started with the simplest case, where the delay and the delay rate are entirely due to the VLBI geometry. Then, we have considered all their components. In all the cases, the effects of polar motion, UT1-UTC, precession and nutation are neglected.

THE METHOD

In the first approach, we assume that the delay and the delay rate are due to the VLBI geometry only and we write the equations of observations in a geocentric equatorial frame. The baseline vector \vec{B} remains fixed in a terrestrial frame while the unit vector $\vec{\sigma}$ in the direction of the radio source remains fixed in a celestial frame. Be S(O, x, y, z) a geocentric frame fixed in space, with the *x*-axis near the direction of the equinox and the *z*-axis in the direction of the rotation axis of the Earth. S'(O, x', y', z') is a geocentric frame fixed in the Earth, with the x'-axis in the direction of the Greenwich meridian; as we are neglecting the effects of polar motion, precession and nutation, the *z*-axis of both frames are coincident.

The declination δ_B of \vec{B} remains constant in the frame fixed in space, but its right ascension α_B changes with time as the right ascension of the Greenwich meridian does.

The arc A between two sources of coordinates (α_1, δ_1) , (α_2, δ_2) can be written:

$$\cos A = \sin \delta_1 \sin \delta_2 + \cos \delta_1 \cos \delta_2 \cos(\alpha_1 - \alpha_2) \tag{1}$$

For each radio source *i* the technique of VLBI determines the delay τ_i and the delay rate $\dot{\tau}_i$, which can be written in terms of the radio source and the baseline coordinates. The observation of the pair of sources gives four equations. Combining them with eq. (1) the arc A adopts the form:

$$\cos A = f(\tau_1, \dot{\tau}_1, \tau_2, \dot{\tau}_2, \sec^2 \delta_B, \sec^4 \delta_B, B, \partial \alpha_G(t) / \partial t)$$
(2)

We have eliminated the radio sources coordinates $(\alpha_1, \delta_1), (\alpha_2, \delta_2)$ and the right ascension α_B of the baseline, but its declination δ_B is still present. The dependence on the declination of the baseline makes this expression useless. Our goal is to eliminate all the parameters except the fundamental VLBI observables τ and $\dot{\tau}$. It is probable that the equatorial system we have used to write our equations is not the best adapted to our analysis.

As the arc is independent of the system adopted, we can choose a set of arbitrary directions in space. We can define a new orthogonal frame associated to the baseline and to the angular rotation velocity of the Earth $\vec{\Omega}$. The direction of the new axes $\hat{\imath}_1, \hat{\imath}_2, \hat{\imath}_3$ are given by

$$\begin{aligned} \hat{\imath}_1 &= \hat{e} \\ \hat{\imath}_2 &= \hat{\omega} \wedge \hat{\imath}_1 \\ \hat{\imath}_3 \perp \hat{\imath}_2 \perp \hat{\imath}_1 \end{aligned}$$

where \hat{e} and $\hat{\omega}$ are the unit vectors of \vec{B} and $\vec{\Omega}$ respectively.

Although the modules of the vectors \vec{B} , $\vec{\Omega}$ and the arc A are unknown we can adopt a priori values B_0 , Ω_0 , A_0 and include in the unknowns their dimensionless corrections ΔB , $\Delta \Omega$, ΔA .

Writing the observables τ_1 , $\dot{\tau}_1$, τ_2 , $\dot{\tau}_2$ in the new reference system we obtain the differential expression:

$$\sin A_0 \ \Delta A = \cos A_0 + C_1 + C_2 \ \Delta B + C_3 \ \Delta \Omega \tag{3}$$

where we have consider only the first orders in the small corrections ΔB , $\Delta \Omega$, ΔA , as they are $\approx 10^{-8}$. The coefficients C_i in eq (3) depend on the known parameters B_0 , Ω_0 , τ_i , $\dot{\tau}_i$.

Let us consider now a more realistic expression of the delay and the delay rate. If p(t) takes into account the effects of the propagation medium and q_0 , q are the clock constants, they can be written as:

$$\tau = \tau_g + p(t) + q_0 + q \Delta t$$
$$\dot{\tau} = \dot{\tau}_g + \dot{p}(t) + q$$

where τ_g is the geometric component of the delay.

Under these conditions, the correction ΔA to the arc can be expressed as:

 $\sin A_0 \ \Delta A = \cos A_0 + C_1 \ \Delta B + C_2 \ Q_0 + C_3 \ Q + C_4 \ \Delta \Omega + C_5 \ P + C_6 \ \dot{P} + C_7 \quad (4)$

up to the first order in the small quantities ΔB , $\Delta \Omega$, p, q, q_0 , \dot{p} .

The theoretical precission of this method is given by the order of the terms neglected. The largest contribution is due to the term in $\Delta\Omega^2$, that is the correction to the module of the angular rotation velocity of the Earth. Adopting standard values of the known parameters it is aproximately 0.0004"

Observing the pair of sources with a three-element interferometer we can define a reference system independently of the Earth rotation vector. This new system will be attached to the unit vectors \hat{e}_1 , \hat{e}_2 in the directions of two of the baselines in the form:

$$\begin{aligned} \hat{\imath}_1' &= \hat{e}_1 \\ \hat{\imath}_2' &= \hat{e}_1 \wedge \hat{e}_2 \\ \hat{\imath}_3' \perp \hat{\imath}_2' \perp \hat{\imath}_1' \end{aligned}$$

The observed delays are (first subscript denotes baseline and second denotes source):

$$\begin{aligned} \tau_{11} &= -\frac{\vec{B}_1 \cdot \hat{e}_1}{c} + p_{11}(t) + q_{01} + q_1 \ \Delta t \\ \tau_{12} &= -\frac{\vec{B}_1 \cdot \hat{e}_2}{c} + p_{12}(t) + q_{01} + q_1 \ \Delta t \\ \tau_{21} &= -\frac{\vec{B}_2 \cdot \hat{e}_1}{c} + p_{21}(t) + q_{02} + q_2 \ \Delta t \\ \tau_{22} &= -\frac{\vec{B}_2 \cdot \hat{e}_2}{c} + p_{22}(t) + q_{02} + q_2 \ \Delta t \end{aligned}$$

From these last four equations we can obtain an expression for the correction ΔA to the arc wich is analogous to eq (4). A new parameter is addeed, the relative positions of baselines D, wich can be written as $D = D_0(1 + \Delta D)$, being D_0 an adopted value and ΔD an unknown correction of the order 10^{-8} .

We have analysed the particular case in wich the observations are conducted with one interferometer in the East-West direction and the other in the North-South direction $(D_0 = 90^\circ)$. Finally:

$$\Delta A \sin A_0 = \cos A_0 + C_0 + C_1 \ \Delta B_1 + C_2 \ \Delta B_2 + C_3 \ \Delta D + C_4 \ P_{11} + C_5 \ P_{12} + C_6 \ P_{21} + C_7 \ P_{22} + C_8 \ Q_{01} + C_9 \ Q_{02}$$
(5)
+ $C_{10} \ Q_1 \ \Delta T + C_{11} \ Q_2 \ \Delta T$

where $C_i = C_i(B_{01}, B_{02}, D_0, \tau_{11}, \tau_{12}, \tau_{21}, \tau_{22})$.

The largest terms neglected are:

$$\frac{2c^2}{B^2} \tau \Delta B q_0 \approx 10^{-7''}$$
$$\frac{4c^4}{B^4} \tau^3 \Delta B q_0 \approx 4 \times 10^{-8''}$$

CONCLUSION

We have shown that it is possible to obtain a set of equations of observations independently of the baseline orientation and the vector of angular rotation velocity of the Earth. Adopting a reference system attached to the baseline(s) we can construct a new observable, the arc between a pair of radio sources simultaneously observed.

Observing K pairs of sources, eq (5) gives K equations of observation with K unknowns ΔA and three more unknowns ΔB_1 , ΔB_2 , ΔD . It is preferable to parametrize the system with a minimum number of unknowns transforming the K corrections to the arcs ΔA into corrections ($\Delta \phi, \Delta \epsilon$) to arbitrary spherical coordinates of the radio sources.

In this new system the origins are fixed by adopting a priori values for three coordinates i.e. ϕ , ϵ of a source and ϵ of another one. In the new parameters, eq (5) becomes:

$$\sin A_0 \ \Delta A = f_1 \ \Delta \epsilon_1 + f_2 \ \Delta \epsilon_2 + f_3 \ (\Delta \phi_1 - \Delta \phi_2) \tag{6}$$

Five radio sources are enough to obtain a solution. However, we prefer to have a superabundant system to solve it with a least squares fit.

The theoretical precission of the method, estimated from the terms neglected in the developments, is at the level of 0.0004'' when the observations are performed with a two-element interferometer and the reference system is attached to the unique baseline and to the direction of the vector of rotation velocity of the Earth. When the reference system is defined only by two baseline vectors (case of the three-element interferometer), and only the delay τ is considered in the expression of the arc, the theoretical precission is 10^3 times higher.

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