

## CONTACT METRIC THREE-MANIFOLDS WITH CONSTANT SCALAR TORSION

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### Abstract

In this paper we study three-dimensional contact metric manifolds satisfying  $\|\tau\| = \text{constant}$ . The local description, as well as several global results and new examples of such manifolds are given.

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### 1. Introduction

In contact geometry, the tensor field  $\tau = \mathcal{L}_\xi g$ , introduced by Chern and Hamilton [4], and the Jacobi operator  $l = R(\cdot, \xi)\xi$  play a fundamental role. In a  $(2m + 1)$ -dimensional contact metric manifold  $M$ , the function  $\text{Tr } l$  and the scalar torsion  $\|\tau\|$  are related by the relation  $\|\tau\|^2 = 4(2m - \text{Tr } l) \geq 0$  [4, 13]. So the constancy of each  $\text{Tr } l$  and  $\|\tau\|$  implies the constancy of the other. (Thus we will be using the constancy of  $\|\tau\|$  or  $\text{Tr } l$  interchangeably in this paper.) It is well known that there exist a lot of classes of contact metric manifolds with  $\|\tau\| = \text{constant}$ , such as the Sasakian manifolds, the  $K$ -contact manifolds, the tangent sphere bundle equipped with the Sasaki metric of a Riemannian manifold of constant curvature, or more generally the  $(\kappa, \mu)$ -contact manifolds [2], the normal bundle of a maximal dimension integral submanifold of a Sasakian manifold [1, page 189], the homogeneous contact Riemannian three-manifolds [12], the three-dimensional pseudosymmetric of constant type contact metric manifolds which satisfy one more condition [6, 7], and the Jacobi  $(\kappa, \mu)$ -contact manifolds [5]. For more information about contact metric manifolds with  $\|\tau\| = \text{constant}$ , see [9, 11]. So it is natural to look for the potential existence of more contact metric manifolds with  $\|\tau\| = \text{constant}$  beyond the aforementioned well-known classes.

In this paper, we study the condition  $\|\tau\| = \text{constant}$  in the three-dimensional case and the content is organized in the following way. Section 2 is devoted mainly to

preliminaries on contact metric manifolds and to some new examples. In Section 3, several global results of three-dimensional contact metric manifolds with  $\|\tau\| = \text{constant}$  are given. Finally, Section 4 is concerned with the local description of such manifolds. In particular, in this section, in terms of contact metric manifolds with  $\|\tau\| = \text{constant}$ , we distinguish between and characterize the  $(\kappa, \mu)$ -contact manifolds and the Jacobi  $(\kappa, \mu)$ -contact manifolds.

### 2. Preliminaries

A contact manifold is a differentiable manifold  $M^{2m+1}$  together with a global 1-form  $\eta$  (a contact form) such that  $\eta \wedge (d\eta)^m \neq 0$  everywhere. Since  $d\eta$  is of rank  $2m$ , there exists a unique vector field  $\xi$  (the Reeb or the characteristic vector field of the contact structure  $\eta$ ) satisfying  $\eta(\xi) = 1$  and  $d\eta(X, \xi) = 0$  for all vector fields  $X$ . The distribution  $D$  defined by the subspace  $X \in T_pM : \eta(X) = 0$  for all  $p \in M$  is called the contact distribution. Every contact manifold has an underlying almost contact structure  $(\eta, \xi, \phi)$ , where  $\phi$  is a global tensor field of type  $(1, 1)$  such that

$$\eta(\xi) = 1, \quad \phi(\xi) = 0, \quad \eta \circ \phi = 0, \quad \phi^2 = -I + \eta \otimes \xi. \tag{2.1}$$

A Riemannian metric  $g$  (the associated metric) can be defined such that

$$\eta(X) = g(X, \xi) \quad \text{and} \quad d\eta(X, Y) = g(X, \phi Y) \tag{2.2}$$

for all vector fields  $X$  and  $Y$  on  $M^{2m+1}$ . We note that  $g$  and  $\phi$  are not unique for a given contact form  $\eta$ , but  $g$  and  $\phi$  are canonically related to each other by

$$g(\phi X, \phi Y) = g(X, Y) - \eta(X)\eta(Y). \tag{2.3}$$

We refer to  $(\eta, \xi, \phi, g)$  as a contact metric structure (c.m.s. in short) and to the manifold  $M^{2m+1}$  carrying such a structure as a contact metric manifold (c.m.m. in short) and this will be denoted by  $M^{2m+1}(\eta, \xi, \phi, g)$ . Denoting Lie differentiation and the curvature tensor by  $\mathcal{L}$  and  $R$ , respectively, we define the operators  $l, h$  and  $\tau$  by

$$l = R(\cdot, \xi)\xi, \quad h = \frac{1}{2}\mathcal{L}_\xi\phi, \quad \tau = \mathcal{L}_\xi g = 2g(h\phi, \cdot).$$

On every c.m.m.  $M^{2m+1}(\eta, \xi, \phi, g)$  we have many important formulas,

$$l\xi = h\xi = 0, \quad \eta \circ h = 0, \quad \text{Tr } h = \text{Tr } \phi h = 0, \quad h\phi = -\phi h, \\ hX = \lambda X \quad \text{implies } h\phi X = -\lambda\phi X.$$

Moreover, if  $\nabla$  is the Riemannian connection of  $g$ ,  $S$  is the Ricci tensor of type  $(0, 2)$ ,  $Q$  is the corresponding Ricci operator satisfying  $g(QX, Y) = S(X, Y)$  and  $r = \text{Tr } Q$  is the scalar curvature, then

$$\nabla_\xi\phi = 0, \quad \nabla_X\xi = -\phi X - \phi hX, \quad \text{Tr } l = g(Q\xi, \xi) = 2m - \text{Tr } h^2 \leq 2m \\ \tau = 2g(\phi \cdot, h \cdot), \quad \nabla_\xi\tau = 2g(\phi \cdot, (\nabla_\xi h) \cdot) \\ \nabla_\xi h = \phi - \phi l - \phi h^2.$$

The conditions  $\|\tau\| = \text{constant}$ ,  $\text{Tr } l = \text{constant}$  and  $\text{Tr } h^2 = \text{constant}$  are equivalent. A c.m.m.  $M^{2m+1}(\eta, \xi, \phi, g)$  for which  $\xi$  is a Killing vector field, that is, for which  $\mathcal{L}_\xi g = 0$ , is called a  $K$ -contact manifold. A c.m.m.  $M^{2m+1}(\eta, \xi, \phi, g)$  is  $K$ -contact manifold if and only if  $h = 0$  (or, equivalently,  $\tau = 0$ ). If we take the product  $M^{2m+1} \times \mathbb{R}$ , the c.m.s. on  $M^{2m+1}$  gives rise to an almost complex structure  $J$  on  $M^{2m+1} \times \mathbb{R}$  given by  $J(X, f(d/dt)) = (\phi X - f\xi, \eta(X)(d/dt))$ . If this structure is integrable, then  $M^{2m+1}$  is called Sasakian. A c.m.m. is Sasakian if and only if  $R(X, Y)\xi = \eta(Y)X - \eta(X)Y$  for all vector fields  $X, Y$  on the manifold. If  $\dim M^{2m+1} = 3$ , then a  $K$ -contact manifold is Sasakian. A c.m.m.  $M(\eta, \xi, \phi, g)$  is said  $H$ -contact manifold if the characteristic vector field  $\xi$  is harmonic or, equivalently, if  $\xi$  is an eigenvector of the Ricci operator [13]. Sasakian and  $K$ -contact manifolds are  $H$ -contact manifolds. More details on contact manifolds are found in [1].

A generalization of Sasakian manifolds are the  $(\kappa, \mu)$ -contact manifolds [2], the curvature tensor of which satisfies the condition

$$R(X, Y)\xi = \kappa(\eta(Y)X - \eta(X)Y) + \mu(\eta(Y)hX - \eta(X)hY) \tag{2.4}$$

for all vector fields  $X, Y$ , where  $\kappa = \text{Tr } l/2$  and  $\mu$  are constant. If  $\kappa, \mu$  in (2.4) are nonconstant smooth functions, then  $M^{2m+1}$  is called a generalized  $(\kappa, \mu)$ -contact manifold [10].

Moreover generalizations of  $(\kappa, \mu)$ -contact manifolds and  $K$ -contact manifolds are the Jacobi  $(\kappa, \mu)$ -contact manifolds, which satisfy the condition

$$l = -\kappa\phi^2 + \mu h, \tag{2.5}$$

where  $\kappa, \mu$  are constant [5]. From (2.5),  $\text{Tr } h = 0$  and  $\text{Tr } \phi^2 = -2m$ , it follows that  $\text{Tr } l = 2m\kappa = \text{constant}$ .

We note that all manifolds are assumed to be connected and smooth. The set of the vector fields on the manifold  $M$  will be denoted by  $\mathcal{X}(M)$ .

In the next proposition, an essential characteristic of the class of contact metric manifolds with  $\text{Tr } l = \text{constant}$  is proved.

**PROPOSITION 2.1.** *For a contact metric  $(2m + 1)$ -manifold, the condition  $\text{Tr } l = \text{constant}$  is invariant under a  $D$ -homothetic deformation.*

**PROOF.** By a  $D$ -homothetic deformation [14] on  $M(\eta, \xi, \phi, g)$  we mean a change of structure tensors of the form

$$\bar{\eta} = \alpha\eta, \quad \bar{\xi} = \frac{1}{\alpha}\xi, \quad \bar{\phi} = \phi, \quad \bar{g} = \alpha g + \alpha(\alpha - 1)\eta \otimes \eta,$$

where  $\alpha$  is a positive constant. It is well known that  $M(\bar{\eta}, \bar{\xi}, \bar{\phi}, \bar{g})$  is also a c.m.m. By direct computation, we see that the tensor  $h$  is transformed in the following way.

$$\bar{h} = \frac{1}{\alpha}h.$$

Moreover, using this and  $\text{Tr } l = \text{constant}$  (equivalently,  $\text{Tr } h^2 = \text{constant}$ ), we get  $\text{Tr } \bar{h}^2 = \text{constant}$  and so  $\text{Tr } \bar{l} = \text{constant}$  for any positive number  $\alpha$ .

In the following, we give new examples of contact metric manifolds with  $\text{Tr } l = \text{constant} \neq 2$ . Examples (1)–(3) concern Jacobi  $(\kappa, \mu)$ -contact manifolds, which, for appropriate choices of the function  $f = f(y, z)$ , degenerate into  $(\kappa, \mu)$ -contact manifolds. Example (4) concerns a Jacobi  $(\kappa, \mu)$ -contact manifold. Examples (5) and (6) concern contact metric manifolds with  $\text{Tr } l = \text{constant}$ , which are not Jacobi  $(\kappa, \mu)$ -contact manifolds.  $\square$

**Examples.** In all six examples, the three-dimensional manifold  $M$  is always the same contact manifold  $(\mathbb{R}^3, \eta = dx - y dz)$ , and only the associated metric  $g$  defines the different examples.

(1) Consider on  $M$  an arbitrary smooth function  $f = f(y, z)$  of variables  $y, z$ . The tensor fields  $(\eta, \xi, \phi, g)$ , where

$$\xi = \frac{\partial}{\partial x},$$

$$g = (g_{ij}) : g_{11} = g_{22} = 1, \quad g_{12} = g_{21} = 0, \quad g_{13} = g_{31} = -y,$$

$$g_{23} = g_{32} = \frac{1}{2}(\rho x - f), \quad g_{33} = y^2 + \frac{1 + (\rho x - f)^2}{4}, \quad \rho = \text{constant} > 0$$

and

$$\phi = (\phi_{ij}) : \phi_{11} = \phi_{21} = \phi_{31} = 0, \quad \phi_{12} = 2y, \quad \phi_{13} = y(\rho x - f), \quad \phi_{22} = f - \rho x,$$

$$\phi_{23} = -\frac{1 + (\rho x - f)^2}{2}, \quad \phi_{32} = 2, \quad \phi_{33} = \rho x - f,$$

define a contact metric structure on  $M$ . Moreover,  $M(\eta, \xi, \phi, g)$  is generally a non $(\kappa, \mu)$ -contact manifold, Jacobi  $(\kappa, \mu)$ -contact manifold with  $\kappa = \text{Tr } l/2 = 1 - \rho^2/4 < 1$  and  $\mu = 2 - \rho < 2$ . In particular, if we choose  $f = f(z)$ , then  $M(\eta, \xi, \phi, g)$  is a  $(\kappa, \mu)$ -contact manifold.

(2) Consider on  $M$  an arbitrary smooth function  $f = f(y, z)$  of variables  $y, z$ . The tensor fields  $(\eta, \xi, \phi, g)$ , where

$$\xi = \frac{\partial}{\partial x},$$

$$g = (g_{ij}) : g_{11} = 1, \quad g_{12} = g_{21} = 0, \quad g_{13} = g_{31} = -y, \quad g_{22} = e^{\rho x}$$

$$g_{23} = g_{32} = -\frac{1}{2}f e^{\rho x}, \quad g_{33} = y^2 + \frac{1 + f^2 e^{2\rho x}}{4e^{\rho x}}, \quad \rho = \text{constant} > 0$$

and

$$\phi = (\phi_{ij}) : \phi_{11} = \phi_{21} = \phi_{31} = 0, \quad \phi_{12} = 2y e^{\rho x}, \quad \phi_{13} = -y f e^{\rho x},$$

$$\phi_{22} = f e^{\rho x}, \quad \phi_{32} = 2e^{\rho x}, \quad \phi_{23} = -\frac{1 + f^2 e^{2\rho x}}{2e^{\rho x}}, \quad \phi_{33} = -f e^{\rho x},$$

define a contact metric structure on  $M$ . Moreover,  $M(\eta, \xi, \phi, g)$  is generally a non $(\kappa, \mu)$ -contact manifold, Jacobi  $(\kappa, \mu)$ -contact manifold with  $\kappa = \text{Tr } l/2 = 1 - \rho^2/4$  and  $\mu = 2$ .

In particular,  $M$  is a  $(\kappa, \mu)$ -contact manifold if we choose  $f(y, z) = -\frac{1}{2}\rho y^2 + d(z)$ , where  $d(z)$  is a smooth function of  $z$ .

(3) Consider on  $M$  an arbitrary smooth function  $f = f(y, z)$  of variables  $y, z$ . The tensor fields  $(\eta, \xi, \phi, g)$ , where

$$\xi = \frac{\partial}{\partial x},$$

$$g = (g_{ij}) : g_{11} = 1, \quad g_{12} = g_{21} = 0, \quad g_{13} = g_{31} = -y, \quad g_{23} = g_{32} = -\frac{1}{2}(f + \rho x),$$

$$g_{22} = 1, \quad g_{33} = y^2 + \frac{1 + (f + \rho x)^2}{4}, \quad \rho = \text{constant} > 0$$

and

$$\phi = (\phi_{ij}) : \phi_{11} = \phi_{21} = \phi_{31} = 0, \quad \phi_{12} = 2y, \quad \phi_{13} = -y(f + \rho x),$$

$$\phi_{22} = f + \rho x, \quad \phi_{32} = 2, \quad \phi_{23} = -\frac{1 + (f + \rho x)^2}{2}, \quad \phi_{33} = -(f + \rho x)$$

define a contact metric structure on  $M$ . Moreover,  $M(\eta, \xi, \phi, g)$  is generally a non $(\kappa, \mu)$ -contact manifold, Jacobi  $(\kappa, \mu)$ -contact manifold with  $\kappa = \text{Tr } l/2 = 1 - \rho^2/4$  and  $\mu = \rho + 2 > 2$ . In particular,  $M$  is a  $(\kappa, \mu)$ -contact manifold if we choose  $f = f(z)$ .

(4) The tensor fields  $(\eta, \xi, \phi, g)$ , where

$$\xi = \frac{\partial}{\partial x},$$

$$g = (g_{ij}) : g_{11} = 1, \quad g_{12} = g_{21} = 0, \quad g_{13} = g_{31} = -y, \quad g_{23} = g_{32} = \frac{\rho x}{2},$$

$$g_{22} = \rho^2 x^2 + 1, \quad g_{33} = y^2 + \frac{1}{4}, \quad \rho = \text{constant} > 0$$

and

$$\phi = (\phi_{ij}) : \phi_{11} = \phi_{21} = \phi_{31} = 0, \quad \phi_{12} = 2y(\rho^2 x^2 + 1), \quad \phi_{13} = \rho xy,$$

$$\phi_{22} = -\rho x, \quad \phi_{32} = 2(\rho^2 x^2 + 1), \quad \phi_{23} = -\frac{1}{2}, \quad \phi_{33} = \rho x$$

define a contact metric structure on  $M$ . Moreover,  $M(\eta, \xi, \phi, g)$  is a non $(\kappa, \mu)$ -contact manifold, Jacobi  $(\kappa, \mu)$ -contact manifold, with  $\kappa = \text{Tr } l/2 = 1 - \rho^2/4$  and  $\mu = \rho + 2 > 2$ .

(5) In the open subset  $U = \{(x, y, z) \in \mathbb{R}^3 : 0 < y < \pi\}$  of  $M$ , the tensor fields  $(\eta, \xi, \phi, g)$ , where

$$\xi = \frac{\partial}{\partial x},$$

$$g = (g_{ij}) : g_{11} = 1, \quad g_{12} = g_{21} = 0, \quad g_{13} = g_{31} = -y, \quad g_{22} = e^{\rho x \sin y}$$

$$g_{23} = g_{32} = \frac{1}{2} \cot y, \quad g_{33} = y^2 + \frac{1 + \cot^2 y}{4e^{\rho x \sin y}}, \quad \rho = \text{constant} > 0$$

and

$$\begin{aligned} \phi = (\phi_{ij}) : \phi_{11} = \phi_{21} = \phi_{31} = 0, \quad \phi_{12} = 2ye^{\rho x \sin y}, \quad \phi_{13} = y \cot y, \\ \phi_{22} = -\cot y, \quad \phi_{23} = -\frac{1 + \cot^2 y}{2e^{\rho(\sin y)x}}, \quad \phi_{32} = 2e^{\rho x \sin y}, \quad \phi_{33} = \cot y \end{aligned}$$

define a contact metric structure. Moreover,  $U(\eta, \xi, \phi, g)$  is a non-Jacobi  $(\kappa, \mu)$ -manifold, contact metric manifold with  $\text{Tr } l = 2(1 - \rho^2/4) = \text{constant}$  and  $\mu = 2 + \rho \cos y$  is the nonconstant smooth function of Proposition 3.1. This follows a comparison of the Lie brackets  $[\xi, e]$  of Lemma 3.2 and Theorem 4.4 and using  $\mu = -2A, \lambda = \rho/2$ .

(6) Consider on  $M$  the function  $F = \int e^{\rho \cos x} \cos x \, dx, \rho = \text{constant} > 0$ . The tensor fields  $(\eta, \xi, \phi, g)$ , where

$$\begin{aligned} \xi = \frac{\partial}{\partial x}, \\ g = (g_{ij}) : g_{11} = 1, \quad g_{12} = g_{21} = 0, \quad g_{13} = g_{31} = -y, \quad g_{22} = e^{-\rho \cos x} \\ g_{23} = g_{32} = -\frac{F}{2}\rho e^{-\rho \cos x}, \quad g_{33} = y^2 + \frac{1 + \rho^2 F^2 e^{-2\rho \cos x}}{4e^{-\rho \cos x}} \end{aligned}$$

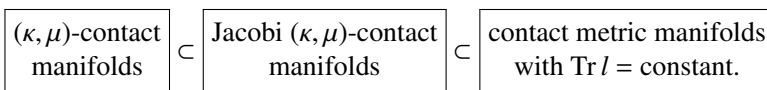
and

$$\begin{aligned} \phi = (\phi_{ij}) : \phi_{11} = \phi_{21} = \phi_{31} = 0, \quad \phi_{12} = 2ye^{-\rho \cos x}, \quad \phi_{13} = -\rho ye^{-\rho \cos x} F, \\ \phi_{22} = \rho e^{-\rho \cos x} F, \quad \phi_{23} = -\frac{1 + \rho^2 F^2 e^{-2\rho \cos x}}{2e^{-\rho \cos x}}, \quad \phi_{32} = 2e^{-\rho \cos x}, \\ \phi_{33} = -\rho e^{-\rho \cos x} F, \end{aligned}$$

define a contact metric structure on  $M$ . Moreover,  $M(\eta, \xi, \phi, g)$  is a non-Jacobi  $(\kappa, \mu)$ -manifold, contact metric manifold with  $\text{Tr } l = 2(1 - \rho^2/4) = \text{constant}$  and  $\mu = 1 + \rho \cos x$  is the nonconstant smooth function of Proposition 3.1. This follows a comparison of the Lie brackets  $[\xi, e]$  of Lemma 3.2 and Theorem 4.4 and using  $\mu = -2A, \lambda = \rho/2$ .

The claims of examples 1–6 could follow from Theorems 4.2, 4.5 and 4.6 by properly choosing the functions  $t = t(x, y, z), c_1 = c_1(y, z)$  and  $c_2 = c_2(y, z)$ . Specifically, examples 1, 2 and 3 follow choosing  $t = \pi, \pi/2, 0$ , respectively, and  $c_1 = f(y, z), c_2 = 2$ . Example 4 follows choosing  $t = 2 \cot^{-1} \rho x, c_1 = 0, c_2 = 2$ . Examples 5 and 6 follow choosing  $t = y, t = x$ , respectively, and  $c_1 = 0, c_2 = 2$ .

From the above examples, it follows that the class of contact metric manifolds with  $\text{Tr } l = \text{constant}$  is a proper generalization of classes of  $(\kappa, \mu)$ -contact manifolds and Jacobi  $(\kappa, \mu)$ -contact manifolds. In particular, the following diagram is valid.



### 3. Global results

As we have seen, the  $(\kappa, \mu)$ -contact manifolds are characterized by the relation (2.4), where  $\kappa = \text{Tr } l/2$  and  $\mu$  are constant. In the next proposition, an expression of  $R(X, Y)\xi$  is given for an arbitrary three-dimensional c.m.m. with  $\text{Tr } l = \text{constant}$ .

**PROPOSITION 3.1.** *Let  $M(\eta, \xi, \phi, g)$  be a three-dimensional c.m.m. with  $\text{Tr } l = \text{constant}$ .*

- (i) *If  $\text{Tr } l = 2$ , then  $M$  is a Sasakian manifold and so  $R(X, Y)\xi = \eta(Y)X - \eta(X)Y$ .*
- (ii) *If  $\text{Tr } l \neq 2$ , then*

$$R(X, Y)\xi = g(X, \phi Y)\phi Q\xi + \kappa(\eta(Y)X - \eta(X)Y) + \mu(\eta(Y)hX - \eta(X)hY)$$

for any  $X, Y \in \mathcal{X}(M)$ , where  $\kappa = \text{Tr } l/2$ , and

$$\mu = -\frac{1}{2}\left(r - 2\kappa - \frac{1}{1 - \kappa} \text{div } \phi h Q\xi\right)$$

is a smooth function, not necessarily constant (compare with examples 5 and 6) where  $\text{div}$  denotes the divergence.

To prove Proposition 3.1, we will need the following lemma [3, 8].

**LEMMA 3.2.** *Let  $M(\eta, \xi, \phi, g)$  be a three-dimensional c.m.m. and let  $U$  be the open set of  $M$ , where  $h \neq 0$ . Then, for any point  $P \in U$ , there exists a smooth local orthonormal basis  $\{\xi, e, \phi e\}$ , such that  $he = \lambda e$ ,  $h\phi e = -\lambda\phi e$ , where  $\lambda$  is a nonvanishing smooth function. Therefore, in  $U$ ,*

$$\left. \begin{aligned} \nabla_e \xi &= -(1 + \lambda)\phi e, & \nabla_{\phi e} \xi &= (1 - \lambda)e, & \nabla_{\xi} \xi &= 0, \\ \nabla_{\xi} e &= A\phi e, & \nabla_e e &= B\phi e, & \nabla_{\phi e} \phi e &= Ce, & \nabla_{\xi} \phi e &= -Ae, \\ \nabla_{\phi e} e &= -C\phi e + (\lambda - 1)\xi, & \nabla_e \phi e &= -Be + (1 + \lambda)\xi, \\ [\xi, e] &= (A + \lambda + 1)\phi e, & [\xi, \phi e] &= -(A - \lambda + 1)e, \\ [e, \phi e] &= -Be + C\phi e + 2\xi, \end{aligned} \right\} \tag{3.1}$$

where  $A, B, C$  are smooth functions on  $U$ . Moreover,

$$\left. \begin{aligned} R(e, \phi e)\xi &= (2\lambda C - e\lambda)e + (\phi e\lambda - 2\lambda B)\phi e, \\ R(e, \xi)\xi &= (1 - \lambda^2 - 2\lambda A)e, & R(\phi e, \xi)\xi &= (1 - \lambda^2 + 2\lambda A)\phi e, \end{aligned} \right\} \tag{3.2}$$

$$\left. \begin{aligned} Qe &= \left(\frac{r}{2} + \lambda^2 - 1 - 2\lambda A\right)e + (\xi\lambda)\phi e + (2\lambda B - \phi e\lambda)\xi, \\ Q\phi e &= (\xi\lambda)e + \left(\frac{r}{2} + \lambda^2 - 1 + 2\lambda A\right)\phi e + (2\lambda C - e\lambda)\xi, \\ Q\xi &= (2\lambda B - \phi e\lambda)e + (2\lambda C - e\lambda)\phi e + (\text{Tr } l)\xi, \end{aligned} \right\} \tag{3.3}$$

$$r = 2(eC + \phi eB - B^2 - C^2 + 2A + 1 - \lambda^2), \tag{3.4}$$

$$B = -\text{div } \phi e, \quad C = -\text{div } e, \tag{3.5}$$

$$\left. \begin{aligned} \xi B &= -C(A - \lambda + 1) + e(A - \lambda), \\ \xi C &= B(A + \lambda + 1) - \phi e(A + \lambda). \end{aligned} \right\} \tag{3.6}$$

**REMARK 3.3.** When  $\text{Tr } l = \text{constant} \neq 2$ , then, from the relation  $h^2 = (\text{Tr } l/2 - 1)\phi^2$ , which is valid on any three-dimensional c.m.m., we have  $h \neq 0$  (that is,  $\lambda \neq 0$ ) in any point of the manifold and so Lemma 3.2 is applied around any point of the manifold. We suppose that  $\lambda > 0$ .

**PROOF OF PROPOSITION 3.1.** (i) If  $\text{Tr } l = 2$ , then  $h = 0$  and so  $M$  is a Sasakian manifold.  
(ii) If  $\text{Tr } l \neq 2$ , let

$$X = p_1e + p_2\phi e + \eta(X)\xi \quad \text{and} \quad Y = \mu_1e + \mu_2\phi e + \eta(Y)\xi \quad \text{for all } X, Y \in X(M),$$

where  $p_i, \mu_i$  are smooth functions on the manifold. Using the basic properties of the curvature tensor and (3.2), we calculate

$$\begin{aligned} R(X, Y)\xi &= (p_1\mu_2 - p_2\mu_1)R(e, \phi e)\xi + (p_1\eta(Y) - \mu_1\eta(X))R(e, \xi)\xi \\ &\quad + (p_2\eta(Y) - \mu_2\eta(X))R(\phi e, \xi)\xi \\ &= (p_1\mu_2 - p_2\mu_1)((2\lambda C)e - (2\lambda B)\phi e) \\ &\quad + (p_1\eta(Y) - \mu_1\eta(X))(1 - \lambda^2 - 2\lambda A)e \\ &\quad + (p_2\eta(Y) - \mu_2\eta(X))(1 - \lambda^2 + 2\lambda A)\phi e \\ &= (p_1\mu_2 - p_2\mu_1)((2\lambda C)e - (2\lambda B)\phi e) \\ &\quad + (1 - \lambda^2)\{\eta(Y)(p_1e + p_2\phi e) - \eta(X)(\mu_1e + \mu_2\phi e)\} \\ &\quad - 2\lambda A\{\eta(Y)(p_1e - p_2\phi e) - \eta(X)(\mu_1e - \mu_2\phi e)\}. \end{aligned}$$

Now, using the relations  $\phi\xi = 0$ ,  $h\xi = 0$ ,  $he = \lambda e$ ,  $h\phi e = -\lambda\phi e$ ,  $\phi^2e = -e$  and (3.3), we find

$$\begin{aligned} g(X, \phi Y) &= g(p_1e + p_2\phi e + \eta(X)\xi, \mu_1\phi e - \mu_2e) = -(p_1\mu_2 - p_2\mu_1) \\ \phi Q\xi &= \phi((2\lambda B)e + (2\lambda C)\phi e + (\text{Tr } l)\xi) = -((2\lambda C)e - (2\lambda B)\phi e) \\ hX &= \lambda(p_1e - p_2\phi e), \quad hY = \lambda(\mu_1e - \mu_2\phi e). \end{aligned}$$

Substituting the above and  $\text{Tr } l = 2(1 - \lambda^2) = 2\kappa$  in  $R(X, Y)\xi$  gives

$$\begin{aligned} R(X, Y)\xi &= g(X, \phi Y)\phi Q\xi \\ &\quad + (1 - \lambda^2)\{\eta(Y)(X - \eta(X)\xi) - \eta(X)(Y - \eta(Y)\xi)\} \\ &\quad - 2\lambda A\left\{\eta(Y)\frac{1}{\lambda}hX - \eta(X)\frac{1}{\lambda}hY\right\} \\ &= g(X, \phi Y)\phi Q\xi + \kappa(\eta(Y)X - \eta(X)Y) + \mu(\eta(Y)hX - \eta(X)hY), \end{aligned}$$

where  $\mu = -2A$ .

Moreover, using  $\phi h = -h\phi$ , (3.3), (3.5), (3.4) and  $\lambda^2 = 1 - \text{Tr } l/2 = 1 - \kappa$ , we calculate

$$\begin{aligned} \phi h Q\xi &= -h\phi Q\xi = -h\{-(2\lambda C)e + (2\lambda B)\phi e\} = 2\lambda^2(Ce + B\phi e) \\ \text{div } \phi h Q\xi &= 2\lambda^2 \text{div}(Ce + B\phi e) = 2\lambda^2(C \text{div } e + eC + B \text{div } \phi e + \phi eB) \\ &= 2\lambda^2(-C^2 + eC - B^2 + \phi eB) = 2\lambda^2\left(\frac{r}{2} - 2A - (1 - \lambda^2)\right). \end{aligned}$$



So  $(1/2\lambda^2)\operatorname{div} \phi h Q \xi = r/2 + \mu - \kappa$  and thus

$$\mu = -\frac{1}{2}\left(r - 2\kappa - \frac{1}{1 - \kappa}\operatorname{div} \phi h Q \xi\right).$$

This completes the proof of Proposition 3.1. □

An immediate consequence of Proposition 3.1 and the definition of a Jacobi  $(\kappa, \mu)$ -contact manifold (see (2.5)) is the following corollary.

**COROLLARY 3.4.** *A three-dimensional c.m.m.  $M(\eta, \xi, \phi, g)$  with  $\operatorname{Tr} l = \text{constant} \neq 2$  is a Jacobi  $(\kappa, \mu)$ -contact manifold if and only if the function  $r - (1/(1 - \kappa))\operatorname{div} \phi h Q \xi$  is constant. In this case,  $\kappa = \operatorname{Tr} l/2$  and  $\mu = -\frac{1}{2}(r - 2\kappa - (1/(1 - \kappa))\operatorname{div} \phi h Q \xi)$ .*

Another immediate consequence of Proposition 3.1 and of the divergence theorem is the following theorem.

**THEOREM 3.5.** *Let  $M(\eta, \xi, \phi, g)$  be a three-dimensional closed (compact without boundary) c.m.m. with  $\operatorname{Tr} l = \text{constant} \neq 2$ . Then*

$$\int_M (r - 2(\kappa - \mu)) dM = 0.$$

We recall that on any three-dimensional, non-Sasakian,  $(\kappa, \mu)$ -contact manifold,  $r = 2(\kappa - \mu)$  [2] is valid.

Now, in order to prove the next theorem, recall that on any three-dimensional Riemannian manifold the well-known formula

$$\sum_i (\nabla_{e_i} Q) e_i = \frac{1}{2} \operatorname{grad} r \tag{3.7}$$

is valid, where  $e_i, i = 1, 2, 3$  is a local orthonormal frame.

**THEOREM 3.6.** *On any three-dimensional c.m.m. with  $\operatorname{Tr} l = \text{constant} \neq 2$ , the following formula is valid.*

$$\operatorname{div} Q \xi = \xi \frac{r}{2}.$$

In particular, if  $M$  is closed, then

$$\int_M (\xi r) dM = 0.$$

**PROOF.** Using the relations (3.1), (3.3), (3.5) and (3.6) of Lemma 3.2, we calculate

$$\begin{aligned} (\nabla_e Q)e &= \nabla_e Qe - Q\nabla_e e = \nabla_e \left\{ \left( \frac{r}{2} + \lambda^2 - 1 - 2\lambda A \right) e + (2\lambda B)\xi \right\} - Q(B\phi e) \\ &= \left( e \frac{r}{2} - 2\lambda eA \right) e + \left( \frac{r}{2} + \lambda^2 - 1 - 2\lambda A \right) B\phi e + 2\lambda(eB)\xi \\ &\quad - 2\lambda B(1 + \lambda)\phi e - B \left\{ \left( \frac{r}{2} + \lambda^2 - 1 + 2\lambda A \right) \phi e + (2\lambda C)\xi \right\} \\ &= \left( e \frac{r}{2} - 2\lambda eA \right) e - \{ 4\lambda AB + 2\lambda(1 + \lambda)B \} \phi e + 2\lambda(eB - BC)\xi, \end{aligned}$$

$$\begin{aligned}
 (\nabla_{\phi e} Q)\phi e &= \nabla_{\phi e} Q\phi e - Q\nabla_{\phi e}\phi e \\
 &= \nabla_{\phi e} \left\{ \left( \frac{r}{2} + \lambda^2 - 1 + 2\lambda A \right) \phi e + (2\lambda C)\xi \right\} - CQe \\
 &= \left( \phi e \frac{r}{2} + 2\lambda \phi e A \right) \phi e + \left( \frac{r}{2} + \lambda^2 - 1 + 2\lambda A \right) C\phi e + (2\lambda \phi e C)\xi \\
 &\quad + 2\lambda C(1 - \lambda)e - C \left\{ \left( \frac{r}{2} + \lambda^2 - 1 - 2\lambda A \right) e + (2\lambda B)\xi \right\} \\
 &= (4\lambda AC + 2\lambda(1 - \lambda)C)e + \left( \phi e \frac{r}{2} + 2\lambda \phi e A \right) \phi e \\
 &\quad + (2\lambda \phi e C - 2\lambda BC)\xi, \\
 (\nabla_{\xi} Q)\xi &= \nabla_{\xi} Q\xi - Q\nabla_{\xi}\xi = \nabla_{\xi} \{ (2\lambda B)e + (2\lambda C)\phi e + (\text{Tr } l)\xi \} \\
 &= 2\lambda(\xi B)e + 2\lambda AB\phi e + 2\lambda(\xi C)\phi e - (2\lambda AC)e \\
 &= 2\lambda \{ -C(A - \lambda + 1) + eA - AC \} e \\
 &\quad + 2\lambda \{ AB + B(A + \lambda + 1) - \phi e A \} \phi e.
 \end{aligned}$$

From the above,

$$(\nabla_e Q)e + (\nabla_{\phi e} Q)\phi e + (\nabla_{\xi} Q)\xi = \left( e \frac{r}{2} \right) e + \left( \phi e \frac{r}{2} \right) \phi e + 2\lambda(eB + \phi eC - 2BC)\xi. \tag{3.8}$$

On the other hand (3.7), for  $e_1 = e, e_2 = \phi e, e_3 = \xi$ , is written as

$$(\nabla_e Q)e + (\nabla_{\phi e} Q)\phi e + (\nabla_{\xi} Q)\xi = \frac{1}{2} \{ (er)e + (\phi er)\phi e + (\xi r)\xi \}. \tag{3.9}$$

Comparing (3.8) and (3.9),

$$\xi r = 4\lambda(eB + \phi eC - 2BC). \tag{3.10}$$

Also

$$\begin{aligned}
 \text{div } Q\xi &= \text{div} \{ (2\lambda B)e + (2\lambda C)\phi e + (\text{Tr } l)\xi \} \\
 &= (2\lambda B)\text{div } e + e(2\lambda B) + (2\lambda C)\text{div } \phi e + \phi e(2\lambda C) + (\text{Tr } l)\text{div } \xi \\
 &= -2\lambda BC + 2\lambda eB - 2\lambda BC + 2\lambda \phi eC \\
 &= 2\lambda(eB + \phi eC - 2BC).
 \end{aligned} \tag{3.11}$$

From (3.10) and (3.11),

$$2\text{div } Q\xi = \xi r.$$

Moreover, if  $M$  is closed, then

$$\int_M (\xi r) dM = 2 \int_M (\text{div } Q\xi) dM = 0.$$

This completes the proof of Theorem 3.6. □

Next, we provide two cases when a three-dimensional c.m.m. with  $\text{tr } l = \text{constant} \neq 2$  reduces to a  $(\kappa, \mu)$ -contact manifold.

**PROPOSITION 3.7.** *A three-dimensional c.m.m.  $M(\eta, \xi, \phi, g)$  with  $\text{Tr } l = \text{constant} \neq 2$  is an  $H$ -contact manifold if and only if  $M(\eta, \xi, \phi, g)$  is a  $(\kappa, \mu)$ -contact manifold. In particular,  $M$  is locally isometric to one of the following unimodular Lie groups  $\text{SU}(2)$ ,  $\text{SL}(2, R)$ ,  $E(2)$ ,  $E(1, 1)$  equipped with a left invariant metric.*

**PROOF.** If  $M$  is an  $H$ -contact manifold, then  $\phi Q\xi = 0$  and so, from Proposition 3.1,

$$R(X, Y)\xi = \kappa(\eta(Y)X - \eta(X)Y) + \mu(\eta(Y)hX - \eta(X)hY),$$

where  $\kappa = \text{Tr } l/2 = \text{constant} \neq 2$  and  $\mu$  is a function. This means that  $M$  is a generalized  $(\kappa, \mu)$ -contact manifold. Therefore, from [10, Theorem 3.6], we have that the function  $\mu$  is constant and so  $M$  is a  $(\kappa, \mu)$ -contact manifold. The inverse is an immediate consequence of Proposition 3.1. For the rest of the proof, see [2, Theorem 3].

We note that Proposition 3.7 extends [5, Proposition 1.3]. □

**PROPOSITION 3.8.** *If the Ricci operator  $Q$  of a three-dimensional c.m.m. with  $\text{Tr } l = \text{constant} \neq 2$  is parallel ( $\nabla Q = 0$ ), then  $M$  is flat, that is, a  $(0, 0)$ -contact manifold.*

**PROOF.** At first from (3.7) we get  $r = \text{constant}$ . Moreover, using the formulas (3.1), (3.3) and (3.6) of Lemma 3.2, we calculate

$$\begin{aligned} 0 &= (\nabla_\xi Q)e = \nabla_\xi Qe - Q\nabla_\xi e \\ &= \nabla_\xi \left\{ \left( \frac{r}{2} + \lambda^2 - 1 - 2\lambda A \right) e + (2\lambda B)\xi \right\} - Q(A\phi e) \\ &= -2\lambda(\xi A)e + \left( \frac{r}{2} + \lambda^2 - 1 - 2\lambda A \right) A\phi e + 2\lambda(\xi B)\xi \\ &\quad - A \left\{ \left( \frac{r}{2} + \lambda^2 - 1 + 2\lambda A \right) \phi e + (2\lambda C)\xi \right\} \\ &= -2\lambda(\xi A)e - 4\lambda A^2 \phi e + 2\lambda \{-C(2A - \lambda + 1) + eA\}\xi. \end{aligned}$$

Thus

$$A = 0 \quad \text{and} \quad (1 - \lambda)C = 0. \tag{3.12}$$

Following this method and using (3.12) we get, from  $(\nabla_{\phi e} Q)e = 0$  and  $(\nabla_e Q)\phi e = 0$ , the following relations.

$$\begin{aligned} (1 - \lambda)B &= 0, \quad (1 + \lambda)B = 0, \quad (1 + \lambda)C = 0 \\ (\lambda - 1) \left( \frac{r}{2} + \lambda^2 - 1 \right) + 2\lambda\phi e B + 2\lambda C^2 - 2(\lambda - 1)(1 - \lambda^2) &= 0, \tag{3.13} \\ (\lambda + 1) \left( \frac{r}{2} + \lambda^2 - 1 \right) + 2\lambda e C + 2\lambda B^2 - 2(\lambda + 1)(1 - \lambda^2) &= 0. \end{aligned}$$

So, from (3.12), (3.13), (3.4) and  $\text{Tr } l = 2(1 - \lambda^2)$ , we finally find  $Q = 0$ , and from the well-known formula

$$\begin{aligned} R(X, Y)Z &= g(Y, Z)QX - g(X, Z)QY + g(QY, Z)Y - g(QX, Z)Y \\ &\quad - \frac{r}{2} \{g(Y, Z)X - g(X, Z)Y\}, \end{aligned}$$

which is valid on any three-dimensional Riemannian manifold, we have  $R = 0$ . This implies that the manifold is flat. □

#### 4. The local description

In order to describe locally the three-dimensional contact metric manifolds (c.m.m.) with  $\|\tau\| = \text{constant}$  (equivalently,  $\text{Tr } l = \text{constant}$ ), we will use the following classical theorem of Darboux [1, page 24] for the 3-dimensional case.

**THEOREM 4.1.** *For each point  $P$  of a three-dimensional contact manifold  $(M, \eta)$  there exist local coordinates  $(U, (x, y, z))$ ,  $P \in U$ , such that*

$$\eta = dx - y dz. \quad (4.1)$$

Now, let  $M(\eta, \xi, \phi, g)$  be a three-dimensional c.m.m. Our initial goal is to describe  $(\eta, \xi, \phi, g)$  in this Darboux coordinate system.

We have  $\xi = \partial/\partial x$  and, from (2.1),  $\phi(\partial/\partial x) = 0$ . Let

$$\phi \frac{\partial}{\partial y} = a \frac{\partial}{\partial x} + b \frac{\partial}{\partial y} + c \frac{\partial}{\partial z}, \quad (4.2)$$

where  $a, b, c$  are smooth functions on  $U$ . From (4.1),

$$\eta\left(\frac{\partial}{\partial y}\right) = 0, \quad \eta\left(\frac{\partial}{\partial z}\right) = -y. \quad (4.3)$$

From (2.1), (4.2) and (4.3), it follows that

$$c\phi \frac{\partial}{\partial z} = -ab \frac{\partial}{\partial x} - (1 + b^2) \frac{\partial}{\partial y} - bc \frac{\partial}{\partial z} \quad (4.4)$$

and  $a = cy$ . So (4.2) is written as

$$\phi \frac{\partial}{\partial y} = cy \frac{\partial}{\partial x} + b \frac{\partial}{\partial y} + c \frac{\partial}{\partial z}. \quad (4.5)$$

From (4.5), it immediately follows that  $c \neq 0$  everywhere on  $U$  and so (4.4) is written as

$$\phi \frac{\partial}{\partial z} = -by \frac{\partial}{\partial x} - \frac{1 + b^2}{c} \frac{\partial}{\partial y} - b \frac{\partial}{\partial z}. \quad (4.6)$$

Consequently, the matrix of the components of  $\phi$  in this system is given by

$$\phi = \begin{pmatrix} 0 & yc & -yb \\ 0 & b & -\frac{1 + b^2}{c} \\ 0 & c & -b \end{pmatrix}. \quad (4.7)$$

Now, for the calculation of the metric tensor  $g$ , using (2.1), (2.2), (4.3), (4.5), (4.6) and  $d\eta(X, Y) = \frac{1}{2}(X\eta(Y) - Y\eta(X) - \eta([X, Y]))$  (see [1, page 69]), we finally get

$$g_{11} = g\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial x}\right) = 1, \quad g_{12} = g_{21} = g\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}\right) = 0,$$

$$g_{13} = g_{31} = g\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial z}\right) = -y$$

and

$$bg_{22} + cg_{23} = 0, \quad \frac{1 + b^2}{c}g_{22} + bg_{23} = \frac{1}{2}, \quad bg_{23} + cg_{33} = \frac{1}{2} + cy^2.$$

From the last three equations, we have  $g_{22} = c/2 > 0$ ,  $g_{23} = g_{32} = -b/2$  and  $g_{33} = y^2 + (1 + b^2)/2c$ .

So the matrix of components of  $g$  is

$$g = \begin{pmatrix} 1 & 0 & -y \\ 0 & \frac{c}{2} & -\frac{b}{2} \\ -y & -\frac{b}{2} & y^2 + \frac{1 + b^2}{2c} \end{pmatrix} \quad \text{with } \det g = \frac{1}{4}. \tag{4.8}$$

We will now calculate, in the Darboux coordinates system, the tensor field  $h = \frac{1}{2}\mathcal{L}_\xi\phi$ . Using (4.5),

$$\begin{aligned} 2h \frac{\partial}{\partial y} &= (\mathcal{L}_\xi\phi) \frac{\partial}{\partial y} = \left[ \xi, \phi \frac{\partial}{\partial y} \right] - \phi \left[ \xi, \frac{\partial}{\partial y} \right] \\ &= \left[ \frac{\partial}{\partial x}, cy \frac{\partial}{\partial x} + b \frac{\partial}{\partial y} + c \frac{\partial}{\partial z} \right] - \phi \left[ \frac{\partial}{\partial x}, \frac{\partial}{\partial y} \right] = yc_x \frac{\partial}{\partial x} + b_x \frac{\partial}{\partial y} + c_x \frac{\partial}{\partial z}, \end{aligned}$$

where  $A_x = \partial A / \partial x$ . So

$$2h \frac{\partial}{\partial y} = yc_x \frac{\partial}{\partial x} + b_x \frac{\partial}{\partial y} + c_x \frac{\partial}{\partial z}. \tag{4.9}$$

Analogously, using (4.6),

$$2h \frac{\partial}{\partial z} = -yb_x \frac{\partial}{\partial x} - \left( \frac{1 + b^2}{c} \right)_x \frac{\partial}{\partial y} - b_x \frac{\partial}{\partial z}. \tag{4.10}$$

Consequently, the matrix of  $h$  is

$$h = \begin{pmatrix} 0 & \frac{1}{2}yc_x & -\frac{1}{2}yb_x \\ 0 & \frac{1}{2}b_x & -\frac{1}{2}\left(\frac{1 + b^2}{c}\right)_x \\ 0 & \frac{1}{2}c_x & -\frac{1}{2}b_x \end{pmatrix}. \tag{4.11}$$

From (4.9) and (4.10), it follows that  $h = 0$  if and only if  $b_x = c_x = 0$ . So the metric  $g$  is Sasakian (that is,  $\text{Tr } l = 2$ ) if and only if the functions  $b$  and  $c$  are independent of  $x$  (see [1, page 230]). From now on we suppose that the three-dimensional c.m.m.  $M(\eta, \xi, \phi, g)$  has  $\text{Tr } l = \text{constant} \neq 2$ . From (4.11), we have that the eigenvalues of  $h$  satisfy the equation

$$0 = \begin{vmatrix} -\lambda & \frac{1}{2}yc_x & -\frac{1}{2}yb_x \\ 0 & \frac{1}{2}b_x - \lambda & -\frac{1}{2}\left(\frac{1 + b^2}{c}\right)_x \\ 0 & \frac{1}{2}c_x & -\frac{1}{2}b_x - \lambda \end{vmatrix} = \lambda \left\{ \lambda^2 - \frac{b_x^2}{4} + \frac{1}{4}c_x \left( \frac{1 + b^2}{c} \right)_x \right\}.$$

So, since  $\lambda \neq 0$  (Remark 3.3), it follows that

$$4\lambda^2 = b_x^2 - c_x \left( \frac{1+b^2}{c} \right)_x = \frac{(cb_x - bc_x)^2 + c_x^2}{c^2}. \quad (4.12)$$

We note that at any point of the manifold is  $(b_x, c_x) \neq (0, 0)$ .

Equation (4.12) is written as

$$\left( b_x - b \frac{c_x}{c} \right)^2 + \left( \frac{c_x}{c} \right)^2 = \rho^2, \quad \rho^2 = 4\lambda^2 \quad \text{for all } \rho > 0. \quad (4.13)$$

Putting  $b_x - b(c_x/c) = \rho \cos t$  and  $c_x/c = \rho \sin t$  for any smooth function  $t = t(x, y, z)$ , the differential equation (4.13) is reduced to the system of two differential equations given by

$$\left\{ b_x - b \frac{c_x}{c} - \rho \cos t = 0 \text{ and } c_x - \rho(\sin t)c = 0 \right\}. \quad (4.14)$$

The solutions of this system are

$$\begin{aligned} 0 < c = c(x, y, z) &= c_2(y, z)e^{\rho \int (\sin t) dx}, \\ b = b(x, y, z) &= e^{\rho \int (\sin t) dx} \left\{ c_1(y, z) + \rho \int e^{-\rho \int (\sin t) dx} (\cos t) dx \right\}, \end{aligned} \quad (4.15)$$

where  $c_1(y, z)$  and  $c_2(y, z) > 0$  are arbitrary smooth functions of  $y$  and  $z$ .

So we have proved the following theorem.

**THEOREM 4.2.** *Let  $M(\eta, \xi, \phi, g)$  be a three-dimensional contact metric manifold. Then, around any point of  $M$ , there exist coordinates  $(x, y, z)$  so that the tensor fields  $\eta, \xi, \phi, g$  and  $h$  are given by (4.1),  $\xi = \partial/\partial x$ , (4.7), (4.8) and (4.11), respectively, where  $b = b(x, y, z)$  and  $c = c(x, y, z) > 0$  are arbitrary smooth functions. In particular:*

- (i)  $\text{Tr } l = 2$  (that is,  $M$  is a Sasakian manifold) if and only if the functions  $b$  and  $c$  are independent of  $x$ ; and
- (ii)  $\text{Tr } l = \text{constant} \neq 2$  if and only if the functions  $b$  and  $c$  satisfy (4.15), where  $c_1(y, z)$ ,  $c_2(y, z) > 0$  and  $t(x, y, z)$  are arbitrary smooth functions.

**The eigenvector of  $h$  when  $\text{Tr } l = \text{constant} \neq 2$ .** Let us suppose now that  $X = \rho_1(\partial/\partial x) + \rho_2(\partial/\partial y) + \rho_3(\partial/\partial z)$  is a nonzero eigenvector of  $h$  with  $hX = \lambda X$ ,  $\lambda > 0$ , where  $\rho_i$ ,  $i = 1, 2, 3$  are smooth functions. Then, using (4.9), (4.10) and  $h(\partial/\partial x) = 0$ ,

$$\begin{aligned} hX &= \rho_1 h \frac{\partial}{\partial x} + \rho_2 h \frac{\partial}{\partial y} + \rho_3 h \frac{\partial}{\partial z} \\ &= \frac{1}{2} \rho_2 \left( y c_x \frac{\partial}{\partial x} + b_x \frac{\partial}{\partial y} + c_x \frac{\partial}{\partial z} \right) + \frac{1}{2} \rho_3 \left( -y b_x \frac{\partial}{\partial x} - \left( \frac{1+b^2}{c} \right)_x \frac{\partial}{\partial y} - b_x \frac{\partial}{\partial z} \right). \end{aligned}$$

So

$$2hX = y(\rho_2 c_x - \rho_3 b_x) \frac{\partial}{\partial x} + \left( \rho_2 b_x - \rho_3 \left( \frac{1+b^2}{c} \right)_x \right) \frac{\partial}{\partial y} + (\rho_2 c_x - \rho_3 b_x) \frac{\partial}{\partial z}.$$

From this and from  $2hX = 2\lambda(\rho_1(\partial/\partial x) + \rho_2(\partial/\partial y) + \rho_3(\partial/\partial z))$ , we get the system

$$\left\{ 2\lambda\rho_1 = y(\rho_2c_x - \rho_3b_x), 2\lambda\rho_2 = \rho_2b_x - \rho_3\left(\frac{1+b^2}{c}\right)_x, 2\lambda\rho_3 = \rho_2c_x - \rho_3b_x \right\}.$$

From the first and the third equations of the system, we get  $\rho_1 = y\rho_3$  and so  $X = y\rho_3(\partial/\partial x) + \rho_2(\partial/\partial y) + \rho_3(\partial/\partial z)$ . Hence the above system is reduced to the homogeneous system

$$\left\{ (2\lambda - b_x)\rho_2 + \left(\frac{1+b^2}{c}\right)_x \rho_3 = 0, -c_x\rho_2 + (2\lambda + b_x)\rho_3 = 0 \right\} \tag{4.16}$$

with determinant  $d = 0$ . So, using (4.5) and (4.6), the eigenvectors of  $h$  are  $\xi, X, \phi X$ , where

$$\begin{aligned} \xi &= \frac{\partial}{\partial x}, & X &= y\rho_3 \frac{\partial}{\partial x} + \rho_2 \frac{\partial}{\partial y} + \rho_3 \frac{\partial}{\partial z} \quad \text{and} \\ \phi X &= y(\rho_2c - \rho_3b) \frac{\partial}{\partial x} + \left(\rho_2b - \frac{1+b^2}{c}\rho_3\right) \frac{\partial}{\partial y} + (\rho_2c - \rho_3b) \frac{\partial}{\partial z}, \end{aligned} \tag{4.17}$$

with eigenvalues  $0, \lambda$  and  $-\lambda$ , respectively, where  $\rho_2$  and  $\rho_3$  are solutions of the system (4.16) and  $(b_x, c_x) \neq (0, 0)$  everywhere.

**Special cases.** In this paragraph, we will look for conditions that characterize the  $(\kappa, \mu)$ -contact manifolds and the Jacobi  $(\kappa, \mu)$ -contact manifolds as subclasses of the class of contact metric manifolds with  $\text{Tr } l = \text{constant} \neq 2$ .

First, we state the following lemma, the proof of which immediately follows from relations (4.14).

**LEMMA 4.3.** *Let  $M(\eta, \xi, \phi, g)$  be a three-dimensional contact metric manifold with  $\text{Tr } l = \text{constant} \neq 2$  ( $\Leftrightarrow \lambda = \rho/2 = \text{constant} > 0$ ). Then, at any point  $P$  of the manifold, there exists a neighborhood  $U$  of  $P$  so that at least one of the functions  $b_x + \rho$  and  $b_x - \rho$  does not vanish anywhere in  $U$ . Moreover:*

- if  $b_x + \rho \neq 0$ , then  $\cos t \neq -1$  everywhere in  $U$ ; and*
- if  $b_x - \rho \neq 0$ , then  $\cos t \neq 1$  everywhere in  $U$ .*

Now, we will examine, separately, the cases  $b_x + \rho \neq 0$  everywhere in  $U$  and  $b_x - \rho \neq 0$  everywhere in  $U$ . In each case, we will find at each point of the manifold a local orthonormal frame  $(\xi, e, \phi e)$  of eigenvectors of  $h$ . Next, we will compute the Lie brackets  $[\xi, e], [\xi, \phi e]$  and  $[e, \phi e]$  in order to compare these with the corresponding ones of Lemma 3.2.

**The case  $b_x + \rho \neq 0$  everywhere in  $U$ .** From the second equation of (4.16), we have  $\rho_3 = (c_x/(b_x + \rho))\rho_2$ . Substituting  $\rho_3$  in (4.17) and using (4.14), (4.7) and Lemma 4.3, we calculate

$$\begin{aligned} X &= \frac{yc_x\rho_2}{b_x + \rho} \frac{\partial}{\partial x} + \rho_2 \frac{\partial}{\partial y} + \frac{c_x\rho_2}{b_x + \rho} \frac{\partial}{\partial z} \\ &= \frac{\rho_2}{b_x + \rho} \left\{ yc\rho(\sin t) \frac{\partial}{\partial x} + (\rho + b\rho \sin t + \rho \cos t) \frac{\partial}{\partial y} + c\rho(\sin t) \frac{\partial}{\partial z} \right\} \\ &= \frac{\rho_2\rho}{b_x + \rho} \left\{ (\sin t) \left( yc \frac{\partial}{\partial x} + b \frac{\partial}{\partial y} + c \frac{\partial}{\partial z} \right) + (1 + \cos t) \frac{\partial}{\partial y} \right\}. \end{aligned}$$

Choosing  $\rho_2 = (b_x + \rho)/\rho \neq 0$ , we have the nonzero eigenvectors

$$X = (\sin t)\phi \frac{\partial}{\partial y} + (1 + \cos t) \frac{\partial}{\partial y} \quad \text{and} \quad \phi X = -(\sin t) \frac{\partial}{\partial y} + (1 + \cos t)\phi \frac{\partial}{\partial y}.$$

Moreover, using (2.3), (4.3) and (4.8),

$$\begin{aligned} |X|^2 &= |\phi X|^2 = (\sin^2 t)g\left(\phi \frac{\partial}{\partial y}, \phi \frac{\partial}{\partial y}\right) + (1 + \cos t)^2 g\left(\frac{\partial}{\partial y}, \frac{\partial}{\partial y}\right) \\ &= (\sin^2 t + 1 + \cos^2 t + 2 \cos t)g\left(\frac{\partial}{\partial y}, \frac{\partial}{\partial y}\right) \\ &= c(1 + \cos t). \end{aligned}$$

Hence, the vector fields  $(\xi, e, \phi e)$ , where  $\xi = \partial/\partial x$ ,

$$\begin{aligned} e &= \sqrt{\frac{1}{c(1 + \cos t)}} \left( (1 + \cos t) \frac{\partial}{\partial y} + (\sin t)\phi \frac{\partial}{\partial y} \right) \\ \phi e &= \sqrt{\frac{1}{c(1 + \cos t)}} \left( -(\sin t) \frac{\partial}{\partial y} + (1 + \cos t)\phi \frac{\partial}{\partial y} \right) \end{aligned}$$

define, at any point  $P$  of  $U$ , a smooth local orthonormal frame, such that  $h\xi = 0$ ,  $he = \lambda e$ ,  $h\phi e = -\lambda\phi e$  and  $\lambda = \rho/2 > 0$ . Putting  $E = \sin t/\sqrt{c(1 + \cos t)} \neq 0$ ,  $F = (1 + \cos t)/\sqrt{c(1 + \cos t)} \neq 0$ , we have  $\xi = \partial/\partial x$ ,  $e = F(\partial/\partial y) + E\phi(\partial/\partial y)$  and  $\phi e = -E(\partial/\partial y) + F\phi(\partial/\partial y)$  with  $\begin{vmatrix} F & E \\ -E & F \end{vmatrix} = F^2 + E^2 = 2/c \neq 0$ . Now, using the above expressions of  $\xi, e, \phi e$  and relations (4.7) and (4.14), we will calculate the Lie brackets  $[\xi, e]$ ,  $[\xi, \phi e]$  and  $[e, \phi e]$ .

$$\begin{aligned} [\xi, e] &= \left[ \frac{\partial}{\partial x}, E\phi \frac{\partial}{\partial y} + F \frac{\partial}{\partial y} \right] = E_x\phi \frac{\partial}{\partial y} + F_x \frac{\partial}{\partial y} + E \left[ \frac{\partial}{\partial x}, \phi \frac{\partial}{\partial y} \right] + F \left[ \frac{\partial}{\partial x}, \frac{\partial}{\partial y} \right] \\ &= E_x\phi \frac{\partial}{\partial y} + F_x \frac{\partial}{\partial y} + E \left[ \frac{\partial}{\partial x}, yc \frac{\partial}{\partial x} + b \frac{\partial}{\partial y} + c \frac{\partial}{\partial z} \right] \\ &= E_x\phi \frac{\partial}{\partial y} + F_x \frac{\partial}{\partial y} + E \left( yc_x \frac{\partial}{\partial x} + b_x \frac{\partial}{\partial y} + c_x \frac{\partial}{\partial z} \right) \\ &= E_x\phi \frac{\partial}{\partial y} + F_x \frac{\partial}{\partial y} + E \frac{c_x}{c} \left( yc \frac{\partial}{\partial x} + b \frac{\partial}{\partial y} + c \frac{\partial}{\partial z} - b \frac{\partial}{\partial y} \right) + Eb_x \frac{\partial}{\partial y} \\ &= E_x\phi \frac{\partial}{\partial y} + F_x \frac{\partial}{\partial y} + E \frac{c_x}{c} \phi \frac{\partial}{\partial y} - E \frac{c_x}{c} b \frac{\partial}{\partial y} + Eb_x \frac{\partial}{\partial y} \end{aligned}$$



$$\begin{aligned}
 &= \left(E_x + E \frac{c_x}{c}\right) \phi \frac{\partial}{\partial y} + \left(F_x - E \frac{c_x b}{c} + E b_x\right) \frac{\partial}{\partial y} \\
 &= \left(E_x + E \frac{c_x}{c}\right) \phi \frac{\partial}{\partial y} + \left(F_x + E \frac{c b_x - b c_x}{c}\right) \frac{\partial}{\partial y}. \tag{*}
 \end{aligned}$$

But

$$\begin{aligned}
 E_x + E \frac{c_x}{c} &= \left(\frac{\sin t}{\sqrt{c(1 + \cos t)}}\right)_x + \frac{\rho \sin^2 t}{\sqrt{c(1 + \cos t)}} \\
 &= \left\{t_x(\cos t) \sqrt{c(1 + \cos t)} - \frac{1}{2}(\sin t) \frac{1}{\sqrt{c(1 + \cos t)}}(c_x(1 + \cos t) \right. \\
 &\quad \left. - c(\sin t)t_x)\right\} \frac{1}{c(1 + \cos t)} + \frac{\rho \sin^2 t}{\sqrt{c(1 + \cos t)}} \\
 &= \left\{t_x(\cos t) \sqrt{c(1 + \cos t)} - \frac{1}{2} \frac{\sin t}{\sqrt{c(1 + \cos t)}}(\rho c(\sin t)(1 + \cos t) \right. \\
 &\quad \left. - c(\sin t)t_x)\right\} \frac{1}{c(1 + \cos t)} + \frac{\rho \sin^2 t}{\sqrt{c(1 + \cos t)}} \\
 &= \frac{1}{\sqrt{c(1 + \cos t)}} \left\{\rho \sin^2 t + t_x(\cos t) \right. \\
 &\quad \left. - \frac{1}{2} \frac{\sin t}{c(1 + \cos t)}(\rho c(\sin t)(1 + \cos t) - c(\sin t)t_x)\right\} \\
 &= \frac{1}{\sqrt{c(1 + \cos t)}} \left\{\rho \sin^2 t + t_x(\cos t) - \frac{1}{2} \rho \sin^2 t + \frac{1}{2} \frac{\sin^2 t}{1 + \cos t} t_x\right\} \\
 &= \frac{1}{\sqrt{c(1 + \cos t)}} \left\{\frac{1}{2} \rho \sin^2 t + t_x(\cos t) + \frac{1}{2}(1 - \cos t)t_x\right\} \\
 &= \frac{1}{\sqrt{c(1 + \cos t)}} \left\{\frac{1}{2} \rho \sin^2 t + \frac{1}{2}(\cos t)t_x + \frac{1}{2} t_x\right\} \\
 &= \frac{1}{2} \frac{1}{\sqrt{c(1 + \cos t)}} \left\{\rho(1 - \cos^2 t) + t_x(1 + \cos t)\right\} \\
 &= \frac{1 + \cos t}{2 \sqrt{c(1 + \cos t)}}(\rho - \rho \cos t + t_x) = \frac{1}{2} F(\rho - \rho \cos t + t_x).
 \end{aligned}$$

Also,

$$\begin{aligned}
 F_x + E \frac{c b_x - b c_x}{c} &= F_x + E \rho \cos t = \left(\frac{1 + \cos t}{\sqrt{c(1 + \cos t)}}\right)_x + \rho \frac{\sin t \cos t}{\sqrt{c(1 + \cos t)}} \\
 &= \frac{1}{2} \sqrt{\frac{c}{1 + \cos t}} \frac{-t_x(\sin t)c - (1 + \cos t)c_x}{c^2} + \frac{\rho \sin t \cos t}{\sqrt{c(1 + \cos t)}} \\
 &= \frac{1}{2} \frac{\sin t}{\sqrt{c(1 + \cos t)}} \{-t_x - (1 + \cos t)\rho + 2\rho \cos t\} \\
 &= -E \frac{\rho - \rho \cos t + t_x}{2}.
 \end{aligned}$$

Substituting the two last relations in (\*), we finally get

$$[\xi, e] = \frac{1}{2}(\rho - \rho \cos t + t_x)\phi e. \quad (4.18)$$

Comparing (4.18) with  $[\xi, e] = (A + \lambda + 1)\phi e$  of (3.1), we have  $A + \lambda + 1 = \frac{1}{2}(\rho - \rho \cos t + t_x)$ , where  $\lambda = \rho/2$ , and so  $A = -1 - (\rho/2) \cos t + \frac{1}{2}t_x$ . Substituting  $A$  in  $[\xi, \phi e] = -(A - \lambda + 1)$  of (3.1) gives

$$[\xi, \phi e] = \frac{1}{2}(\rho + \rho \cos t - t_x)e. \quad (4.19)$$

Now, we will compute  $[e, \phi e]$ , using the properties of the Lie bracket.

$$\begin{aligned} [e, \phi e] &= \left[ E\phi \frac{\partial}{\partial y} + F \frac{\partial}{\partial y}, F\phi \frac{\partial}{\partial y} - E \frac{\partial}{\partial y} \right] \\ &= \left[ E\phi \frac{\partial}{\partial y}, F\phi \frac{\partial}{\partial y} \right] - \left[ E\phi \frac{\partial}{\partial y}, E \frac{\partial}{\partial y} \right] + \left[ F \frac{\partial}{\partial y}, F\phi \frac{\partial}{\partial y} \right] - \left[ F \frac{\partial}{\partial y}, E \frac{\partial}{\partial y} \right] \\ &= EF \left[ \phi \frac{\partial}{\partial y}, \phi \frac{\partial}{\partial y} \right] + E\phi \frac{\partial}{\partial y} F\phi \frac{\partial}{\partial y} - F\phi \frac{\partial}{\partial y} E\phi \frac{\partial}{\partial y} \\ &\quad - \left( E^2 \left[ \phi \frac{\partial}{\partial y}, \frac{\partial}{\partial y} \right] + E\phi \frac{\partial}{\partial y} E \frac{\partial}{\partial y} - E \frac{\partial}{\partial y} E\phi \frac{\partial}{\partial y} \right) \\ &\quad + \left( F^2 \left[ \frac{\partial}{\partial y}, \phi \frac{\partial}{\partial y} \right] + F \frac{\partial}{\partial y} F\phi \frac{\partial}{\partial y} - F\phi \frac{\partial}{\partial y} F \frac{\partial}{\partial y} \right) \\ &\quad - \left( EF \left[ \frac{\partial}{\partial y}, \frac{\partial}{\partial y} \right] + F \frac{\partial}{\partial y} E \frac{\partial}{\partial y} - E \frac{\partial}{\partial y} F \frac{\partial}{\partial y} \right) \\ &= \left( E\phi \frac{\partial}{\partial y} F - F\phi \frac{\partial}{\partial y} E \right) \phi \frac{\partial}{\partial y} - \left( F \frac{\partial}{\partial y} E - E \frac{\partial}{\partial y} F \right) \frac{\partial}{\partial y} \\ &\quad - \left( E\phi \frac{\partial}{\partial y} E \frac{\partial}{\partial y} - E \frac{\partial}{\partial y} E\phi \frac{\partial}{\partial y} + E^2 \left[ \phi \frac{\partial}{\partial y}, \frac{\partial}{\partial y} \right] \right) \\ &\quad + \left( F \frac{\partial}{\partial y} F\phi \frac{\partial}{\partial y} - F\phi \frac{\partial}{\partial y} F \frac{\partial}{\partial y} + F^2 \left[ \frac{\partial}{\partial y}, \phi \frac{\partial}{\partial y} \right] \right) \\ &= \left( E\phi \frac{\partial}{\partial y} F - F\phi \frac{\partial}{\partial y} E + E \frac{\partial}{\partial y} E + F \frac{\partial}{\partial y} F \right) \phi \frac{\partial}{\partial y} \\ &\quad + \left( -F \frac{\partial}{\partial y} E + E \frac{\partial}{\partial y} F - E\phi \frac{\partial}{\partial y} E - F\phi \frac{\partial}{\partial y} F \right) \frac{\partial}{\partial y} \\ &\quad + (F^2 + E^2) \left[ \frac{\partial}{\partial y}, \phi \frac{\partial}{\partial y} \right] \\ &= \left( -F^2 \phi \frac{\partial}{\partial y} \frac{E}{F} + \frac{1}{2} \frac{\partial}{\partial y} (E^2 + F^2) \right) \phi \frac{\partial}{\partial y} \\ &\quad + \left( -F^2 \frac{\partial}{\partial y} \frac{E}{F} - \frac{1}{2} \phi \frac{\partial}{\partial y} (E^2 + F^2) \right) \frac{\partial}{\partial y} + (E^2 + F^2) \left[ \frac{\partial}{\partial y}, \phi \frac{\partial}{\partial y} \right]. \quad (**)$$

But using (4.5), we compute

$$\begin{aligned} -F^2 \phi \frac{\partial E}{\partial y F} + \frac{1}{2} \frac{\partial}{\partial y} (E^2 + F^2) &= -\frac{1 + \cos t}{c} \phi \frac{\partial}{\partial y} \frac{\sin t}{1 + \cos t} + \frac{1}{2} \frac{\partial}{\partial y} \frac{2}{c} \\ &= -\frac{1 + \cos t}{c} \frac{1}{1 + \cos t} \phi \frac{\partial}{\partial y} t + \frac{\partial}{\partial y} \frac{1}{c} \\ &= -\frac{1}{c} \left( y c \frac{\partial}{\partial x} + b \frac{\partial}{\partial y} + c \frac{\partial}{\partial z} \right) t - \frac{c_y}{c^2} \\ &= -\frac{1}{c} (y c t_x + b t_y + c t_z) - \frac{c_y}{c^2}. \end{aligned}$$

Also,

$$\begin{aligned} -F^2 \frac{\partial E}{\partial y F} - \frac{1}{2} \phi \frac{\partial}{\partial y} (E^2 + F^2) &= -\frac{1 + \cos t}{c} \frac{\partial}{\partial y} \frac{\sin t}{1 + \cos t} - \frac{1}{2} \phi \frac{\partial}{\partial y} \frac{2}{c} \\ &= -\frac{1 + \cos t}{c} \frac{1}{1 + \cos t} t_y - \left( y c \frac{\partial}{\partial x} + b \frac{\partial}{\partial y} + c \frac{\partial}{\partial z} \right) \frac{1}{c} \\ &= -\frac{1}{c} t_y + \frac{1}{c^2} (y c c_x + b c_y + c c_z) \end{aligned}$$

and

$$\begin{aligned} (E^2 + F^2) \left[ \frac{\partial}{\partial y}, \phi \frac{\partial}{\partial y} \right] &= \frac{2}{c} \left[ \frac{\partial}{\partial y}, y c \frac{\partial}{\partial x} + b \frac{\partial}{\partial y} + c \frac{\partial}{\partial z} \right] \\ &= \frac{2}{c} \left\{ c \frac{\partial}{\partial x} + y c_y \frac{\partial}{\partial x} + b_y \frac{\partial}{\partial y} + c_y \frac{1}{c} \left( \phi \frac{\partial}{\partial y} - y c \frac{\partial}{\partial x} - b \frac{\partial}{\partial y} \right) \right\} \\ &= 2 \frac{\partial}{\partial x} + \frac{2}{c} \left( b_y - \frac{b c_y}{c} \right) \frac{\partial}{\partial y} + \frac{2}{c^2} c_y \phi \frac{\partial}{\partial y} \\ &= 2 \xi + \frac{2}{c^2} (c b_y - b c_y) \frac{\partial}{\partial y} + \frac{2 c_y}{c^2} \phi \frac{\partial}{\partial y}. \end{aligned}$$

Substituting the three last relations in (\*\*\*) gives

$$\begin{aligned} [e, \phi e] &= \left\{ -\frac{1}{c} (y c t_x + b t_y + c t_z) - \frac{c_y}{c^2} \right\} \phi \frac{\partial}{\partial y} + \left\{ -\frac{1}{c} t_y + \frac{1}{c^2} (y c c_x + b c_y + c c_z) \right\} \frac{\partial}{\partial y} \\ &\quad + 2 \xi + \frac{2}{c^2} (c b_y - b c_y) \frac{\partial}{\partial y} + \frac{2 c_y}{c^2} \phi \frac{\partial}{\partial y} \\ &= -\frac{1}{c^2} \{-c_y + c^2 (y t_x + t_z) + b c t_y\} \phi \frac{\partial}{\partial y} \\ &\quad + \frac{1}{c^2} \{-b c_y + c(-t_y + y c_x + c_z + 2 b_y)\} \frac{\partial}{\partial y} + 2 \xi. \end{aligned}$$

So

$$\begin{aligned} [e, \phi e] &= 2 \xi - \frac{1}{c^2} \{-c_y + c^2 (y t_x + t_z) + b c t_y\} \phi \frac{\partial}{\partial y} \\ &\quad + \frac{1}{c^2} \{-b c_y + c(-t_y + y c_x + c_z + 2 b_y)\} \frac{\partial}{\partial y}. \end{aligned} \quad (4.20)$$

**The case  $b_x - \rho \neq 0$  everywhere in  $U$ .** From the first of (4.16), we have  $\rho_2 = (((1 + b^2)/c)_x)/(b_x - \rho)\rho_3$ . Substituting  $\rho_2$  in (4.17) and using (4.14), (4.7) and Lemma 4.3, we finally get

$$X = \frac{\rho_3}{c(1 - \cos t)} \left( (\sin t) \frac{\partial}{\partial y} + (1 - \cos t) \phi \frac{\partial}{\partial y} \right).$$

Choosing  $\rho_3 = c(1 - \cos t) \neq 0$ , we have the nonzero eigenvectors of  $h$

$$X = (\sin t) \frac{\partial}{\partial y} + (1 - \cos t) \phi \frac{\partial}{\partial y} \quad \text{and} \quad \phi X = -(1 - \cos t) \frac{\partial}{\partial y} + (\sin t) \phi \frac{\partial}{\partial y}.$$

Moreover, using (2.3), (4.3) and (4.8),

$$|X| = |\phi X| = \sqrt{c(1 - \cos t)}.$$

Hence, the vector fields  $\xi = \partial/\partial x$ ,

$$e = \frac{1}{\sqrt{c(1 - \cos t)}} \left( (\sin t) \frac{\partial}{\partial y} + (1 - \cos t) \phi \frac{\partial}{\partial y} \right)$$

and

$$\phi e = \frac{1}{\sqrt{c(1 - \cos t)}} \left( -(1 - \cos t) \frac{\partial}{\partial y} + (\sin t) \phi \frac{\partial}{\partial y} \right)$$

define on  $U$  an orthonormal frame of eigenvectors of  $h$ , such that  $h\xi = 0$ ,  $he = \lambda e$ ,  $h\phi e = -\lambda\phi e$  ( $\lambda = \rho/2 > 0$ ). Working as in the case  $b_x + \rho \neq 0$ , we finally get, for the Lie brackets  $[\xi, e]$ ,  $[\xi, \phi e]$ ,  $[e, \phi e]$ , the formulas (4.18), (4.19) and (4.20), respectively. So we have proved the following theorem.

**THEOREM 4.4.** *Let  $M(\eta, \xi, \phi, g)$  be a three-dimensional c.m.m. with  $\text{Tr } l = \text{constant} \neq 2$ . Then, at any point  $P$  of  $M$ , there exists a neighborhood  $U$  of  $P$  so that at least one of the functions  $b_x + \rho$  and  $b_x - \rho$  does not vanish anywhere on  $U$ .*

(i) *If  $b_x + \rho \neq 0$  everywhere in  $U$ , then the triad  $(\xi, e, \phi e)$ , where  $\xi = \partial/\partial x$ ,*

$$e = \frac{1}{\sqrt{(1 + \cos t)}} \left( (1 + \cos t) \frac{\partial}{\partial y} + (\sin t) \phi \frac{\partial}{\partial y} \right)$$

$$\phi e = \frac{1}{\sqrt{c(1 + \cos t)}} \left( -(\sin t) \frac{\partial}{\partial y} + (1 + \cos t) \phi \frac{\partial}{\partial y} \right)$$

*defines a smooth orthonormal frame of eigenvectors of  $h$  in  $U$ , such that  $h\xi = 0$ ,  $he = \lambda e$ ,  $h\phi e = -\lambda\phi e$  ( $\lambda = \rho/2 > 0$ ).*

(ii) *If  $b_x - \rho \neq 0$  everywhere in  $U$ , then the triad  $(\xi, e, \phi e)$ , where  $\xi = \partial/\partial x$ ,*

$$e = \frac{1}{\sqrt{c(1 - \cos t)}} \left( (\sin t) \frac{\partial}{\partial y} + (1 - \cos t) \phi \frac{\partial}{\partial y} \right)$$

$$\phi e = \frac{1}{\sqrt{c(1 - \cos t)}} \left( -(1 - \cos t) \frac{\partial}{\partial y} + (\sin t) \phi \frac{\partial}{\partial y} \right)$$

*defines a smooth orthonormal frame of eigenvectors of  $h$  in  $U$ , such that  $h\xi = 0$ ,  $he = \lambda e$ ,  $h\phi e = -\lambda\phi e$  ( $\lambda = \rho/2 > 0$ ).*

Moreover, in any case ((i) or (ii)), the Lie brackets  $[\xi, e]$ ,  $[\xi, \phi e]$  and  $[e, \phi e]$  are given by

$$\begin{aligned} [\xi, e] &= \frac{1}{2}(\rho - \rho \cos t + t_x)\phi e \\ [\xi, \phi e] &= \frac{1}{2}(\rho + \rho \cos t - t_x)e \\ [e, \phi e] &= 2\xi - \frac{1}{c^2}\{-c_y + c^2(yt_x + t_z) + bct_y\}\phi \frac{\partial}{\partial y} \\ &\quad + \frac{1}{c^2}\{-bc_y + c(-t_y + yc_x + c_z + 2b_y)\}\frac{\partial}{\partial y}. \end{aligned}$$

**When is a three-dimensional c.m.m. with  $\text{Tr } l = \text{constant} \neq 2$  a Jacobi  $(\kappa, \mu)$ -contact manifold?** Comparing relations  $[\xi, e] = \frac{1}{2}(\rho - \rho \cos t + t_x)\phi e$  of Theorem 4.4 and  $[\xi, e] = (A + \lambda + 1)\phi e$  of (3.1), we get  $A + \rho/2 + 1 = \frac{1}{2}(\rho - \rho \cos t + t_x)$  or  $A = \frac{1}{2}(t_x - \rho \cos t) - 1$ . So  $\mu = -2A = \rho \cos t - t_x + 2$ . Hence, using the definition of a Jacobi  $(\kappa, \mu)$ -contact manifold (see relation (2.5)), we state the following theorem.

**THEOREM 4.5.** *Let  $M(\eta, \xi, \phi, g)$  be a three-dimensional c.m.m. with  $\text{Tr } l = \text{constant} \neq 2$ . Then  $M(\eta, \xi, \phi, g)$  is a Jacobi  $(\kappa, \mu)$ -contact manifold if and only if the function  $t = t(x, y, z)$  satisfies the equation*

$$t_x - \rho \cos t + \nu = 0, \quad (4.21)$$

where  $\nu = \text{constant}$ . Therefore, in this case,  $\kappa = \text{Tr } l/2 = 1 - \rho^2/4$  and  $\mu = \nu + 2$ .

Comment. Obviously, the function  $t = t(x, y, z) = \text{constant}$  is a solution of (4.21).

**When is a three-dimensional c.m.m. with  $\text{Tr } l = \text{constant} \neq 2$  a  $(\kappa, \mu)$ -contact manifold?** Let  $M(\eta, \xi, \phi, g)$  be a three-dimensional c.m.m. with  $\text{Tr } l = \text{constant} \neq 2$ . According to Proposition 3.7,  $M$  is a  $(\kappa, \mu)$ -contact manifold if and only if the vector field  $\xi$  is an eigenvector of the Ricci operator  $Q$ , or, equivalently, from (3.3),  $B = C = 0$ , or from (3.1),  $[e, \phi e] = 2\xi$ . Therefore, using  $[e, \phi e] = 2\xi$  and the last relation of Theorem 4.4, we get that  $M(\eta, \xi, \phi, g)$  is a  $(\kappa, \mu)$ -contact manifold if and only if

$$c^2(yt_x + t_z) + bct_y - c_y = 0 \quad \text{and} \quad c(t_y - c_z - 2b_y - yc_x) + bc_y = 0. \quad (4.22)$$

So, we have proved the following theorem.

**THEOREM 4.6.** *Let  $M(\eta, \xi, \phi, g)$  be a three-dimensional c.m.m. with  $\text{Tr } l = \text{constant} \neq 2$ . Then  $M(\eta, \xi, \phi, g)$  is a  $(\kappa, \mu)$ -contact manifold (equivalently, it is  $H$ -contact) if and only if the functions  $b = b(x, y, z)$ ,  $c = c(x, y, z)$  and  $t = t(x, y, z)$  satisfy conditions (4.22).*

It is obvious that, on a  $(\kappa, \mu)$ -contact manifold, the function  $t = t(x, y, z)$  satisfies condition (4.21).

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