Magnetic field of pulsars with superconducting quark core

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Abstract. Within recent nonperturbative approaches to the effective quark interaction the diquark condensate forms a superconductor of second kind. Therefore the magnetic field will not be expelled from the superconducting quark core in accordance with observational data which indicate that life times of pulsar magnetic fields exceed 10^7 years.

The physical properties of pulsars can constrain our hypotheses about the state of matter at high densities. For example, Bailin and Love (1984) have suggested that the magnetic field of pulsars should be expelled from the superconducting interior of the star due to the Meissner effect and decay subsequently within $\approx 10^4$ years. If their arguments would hold in general, the observation of lifetimes as large as 10^7 years (Makashima 1992) would exclude the occurence of an extended superconducting quark matter phase in pulsars. For their estimate, they assumed a homogeneous magnetic field and used a perturbative quark interaction which results in a very small pairing gap. Since both assumptions seem not to be valid in general, we perform a reinvestigation of this question.

Recently there has been excitement (Wilczek 1999) about the observation that in chiral quark models with nonperturbative 4-point interactions the anomalous quark pair amplitudes in the color antitriplet channel can be very large of the order ≈ 100 MeV. In Fig. 1 (a) we show the solution of the corresponding diquark gap equation for a BCS-type quark-quark interaction model (Berges & Rajagopal 1999) and the corresponding Ginzburg-Landau Parameter κ . Quark matter with a diquark condensate appears as a superconductor of second kind into which the magnetic field H can penetrate by forming quantized vortex lines provided $H_{c1} < H < H_{c2}$ where $H_{c1} = 1.8 \cdot 10^{16}$ G (Blaschke, Sedrakian & Shahabasyan 1999).

It is generally accepted that neutrons and protons in the "npe"-phase are superfluid. While the neutrons take part in the rotation, forming a lattice of quantized vortex lines, the superconducting protons will be entrained by the neutrons (Sedrakian & Shahabasian 1980) and form inside the neutron vortex a magnetic field strength H(r) which acts as an external field for the non-entrained protons. This entails the formation of a cluster of proton vortices with fluxes Φ_0 in a region with the radius $\delta_n = 10^{-5}$ cm around the axis of the neutron vortex. The mean magnetic induction within the cluster reaches values of $4 \cdot 10^{14}$ G (Sedrakian & Sedrakian 1995). The magnetic field H(r) generates quark vortex clusters with a radius $\delta_q = 4.3 \cdot 10^{-7}$ cm. Since δ_q is by two orders of magnitude smaller than δ_n , the mean magnetic induction in the clusters of quark vortex lines increases to a value of the order of 10^{18} G, see Fig. 1 (b).



Figure 1. (a) Diquark energy gap Δ and corresponding Ginzburg-Landau parameter κ vs. chemical potential μ (resp. density n) for a BCS-type quark interaction. (b) Magnetic field structure in the interior of a hybrid star with $M = 1.4M_{\odot}$;

The clusters of quark vortex lines which appear due to the entrainment effect in the "npe"-phase will interact with those which are formed by the initial magnetic field (fossil field). Due to this interaction quark vortex lines will not be expelled from the quark core of the star within a time scale of $\tau = 10^4$ years as suggested in (Bailin & Love 1984). This is the basic finding of our paper (Blaschke, Sedrakian & Shahabasyan 1999).

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