

A STATISTICAL ANALYSIS OF THE RADIO PROPERTIES OF A LARGE SAMPLE OF 374 OPTICALLY SELECTED QUASARS

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ABSTRACT. Using a large, unbiased sample of 374 optically selected QSOs, of which 54 are radio detected, we have computed for different redshift ranges the probability distribution function $\psi(R)$ of R defined as the ratio of monochromatic luminosities at 15 GHz and 2500 Å in the QSO's rest frame. A significant variation of $\psi(R)$ with redshift is noticed. At small redshifts ($z < 0.5$), the distinctive feature of $\psi(R)$ is a peak near $R \sim 10^2$, arising due to the QSOs having steep radio spectrum ($\alpha_r < -0.5$).

1. THE QSO SAMPLE AND THE PROBABILITY DISTRIBUTION FUNCTION $\psi(R)$

The sample (374 QSOs) consists of 6 sets, each derived from the radio-surveyed portion of a major QSO list, by excluding all objects with $M_B > -23$ ($q_0=0, H_0=50 \text{ Km s}^{-1} \text{ Mpc}^{-1}$) and also those for which QSO classification or tabulated redshift (z), or app.magnitude (m) has been labelled as doubtful/uncertain in the parent list (Table 1). Radio data available for all 54 detections in the sample permit classification into flat or steep spectrum type, based on α_r defined between ~ 5 and ~ 30 GHz and also yield flux density, S_r , at 15 GHz needed for computing 'R' (see Abstract),

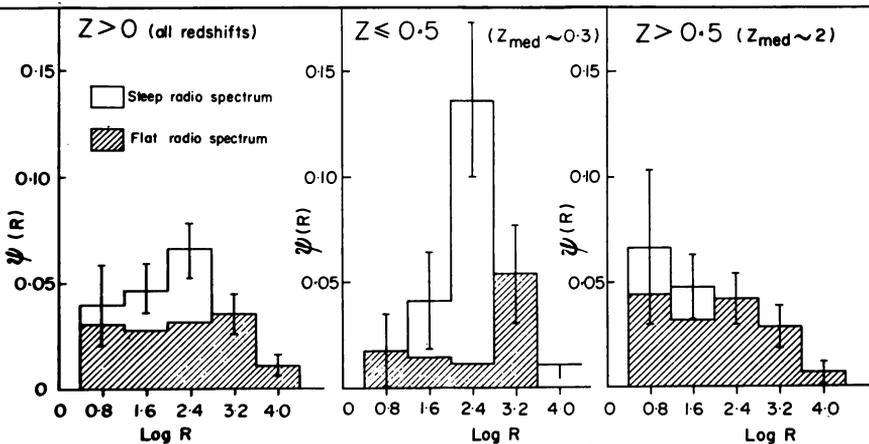


Fig. 1: The probability distribution function $\psi(R)$ for different z -ranges.

Table 1:The 6 sets constituting our sample of 374 optically-selected QSOs

Set, No. of QSOs (radio detections)	Radio telescope used, frequency*, sensitivity limit (m)	Type+ of the adopted mag. and reference to m & z
1. CTIO ¹ : 119(12)	VLA ² , 5 GHz, 10 mJy	m_{con} at $\lambda_e = 1475 \text{ \AA}$, (1)
2. BQS ³ : 92(23)	Effelsberg ⁴ , 10.7 GHz, 10 mJy	B-magnitude, (3)
3. AAO ⁵ : 53(7)	Parkes ⁵ , 5 GHz, 25 mJy	B-magnitude, (5)
4. MCS ⁶ : 48(7)	Effelsberg ⁷ , 5 GHz, 10 mJy	m_{con} at $\lambda_o = 4500 \text{ \AA}$, (6), (ϕ)
5. KP ⁸ : 44(3)	Arecibo ⁹ , 2.4 GHz, ~14 mJy VLA ¹⁰ , 5 GHz, ~1 mJy	m_{con} at $\lambda_o = 4500 \text{ \AA}$, (8), (10)
6. BPS**:	18(2) Westerbork ¹¹ , 1.4GHz, ~4 mJy	B-magnitude, (11)

References (1)Osmer and Smith(Ap.J.Suppl.42,333,1980);(2)Condon et al.(Ap.J.244,5,1981);(3)Schmidt and Green(Ap.J.269,352,1983);(4)Steppe et al.(preprint,1985);(5)Savage and Bolton(MNRAS 188,599, 1979);(6)Lewis et al.(Ap.J.233,787,1979);(7)Strittmatter et al.(Astr.Ap.88,L12,1980);(8)Vaucher and Weedman(Ap.J.240,10,1980);(9)Sramek and Weedman(Ap.J.221,468,1978);(10)Sramek and Weedman (Ap.J.238,435,1980);(11)Marshall et al.(Ap.J.269,42,1983);(12)Fanti et al.(Astr.Ap.61,487,1977).

* flux densities at additional frequencies are taken from the references cited here and from the major radio source catalogues.

+ m_{con} =continuum magnitude, λ_e =emitted wavelength, λ_o =observed wavelength.
 § from the sample given in ref.5.QSOs outside the central $5^\circ \times 5^\circ$ are excluded.
 ϕ values of m not given in ref.6 are deduced following ref.(10).

**derived from the 'Bologna-Palomar Schmidt' sample defined in ref.(11).

following Schmidt and Green (Ap.J.269,352,1983).For the (320) QSOs undetected in radio,the quoted limits to flux density were adopted for S_p . Note that all the 3 frequencies refer to the QSO's rest frame. Now, using this database we computed $\psi(R)$ for the entire z-range (0 to 3.3), following ref.(10) (Fig.1).It is found to be consistent with the earlier determinations(refs.10,12),though much better defined here.We next computed $\psi(R)$ dividing the sample into low ($z < 0.5$) and high($z > 0.5$) redshift subsamples: a clear variation of $\psi(R)$ with z is noticed (Fig.1).At low-z $\psi(R)$ is marked by a peak arising near $R \sim 10^2$ due to steep-spectrum QSOs.

For the high-z subsample, ($z_{med} \sim 2$), $\psi(R)$ is dominated by flat-spectrum QSOs.This is not likely to be due to an observational bias, since for our major sets of both high-z (CTIO,MCS,AAO) and low-z (BQS) QSOs, the survey frequencies roughly correspond to the same emission frequency of ~15 GHz(see Table 1).The inferred lower probability for steep-spectrum QSOs at high z is also difficult to explain as an artefact of the median optical luminosity of our high-z subsample being higher(by ~3 mag.), compared to our low-z subsample. This is because, as seen from fig.2, optically more luminous QSOs in fact seem to produce steep-spectrum sources, preferentially.The inferred rarity of steep-spectrum QSOs at high-z may, thus, be real, unless a bias against sources with steep radio spectrum is somehow present already in the parent optical lists of high-z QSOs(Gopal-Krishna etal,in prep.)

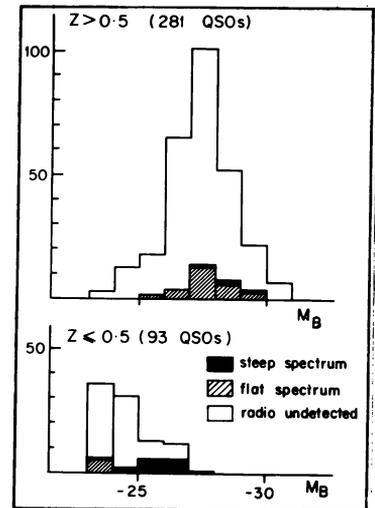


Fig.2:The distribution of M_B for the high-z and low-z QSOs in the sample.